

**MSSM Radiative Corrections**  
**to**  
**Neutrino-nucleon Deep-inelastic Scattering**

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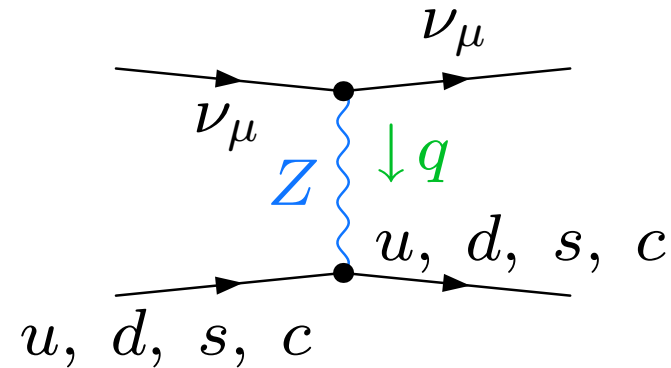
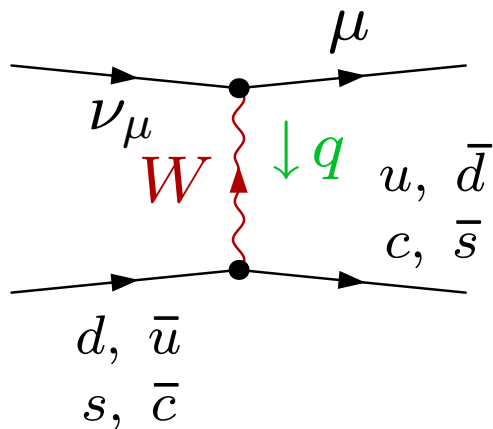
## outline :

- Introduction
  - deep inelastic  $\nu N$  scattering at NuTeV
  - possible explanations
- MSSM radiative corrections to  $\nu_\mu N$  DIS
  - definition of  $\delta R^{\nu, \bar{\nu}} = R_{\text{MSSM}}^{\nu, \bar{\nu}} - R_{\text{SM}}^{\nu, \bar{\nu}}$
  - superpartner loop corrections
  - MSSM-SM Higgs loop difference
- Results for  $\delta R^\nu, \delta R^{\bar{\nu}}$ 
  - how to scan over MSSM parameters?
  - MSSM parameter scan

## • Introduction

– deep inelastic  $\nu N$  scattering at NuTeV

In the SM **neutral** (NC) and **charged** current (CC) neutrino nucleon scattering are described in LO by  $t$ -channel  $W$  and  $Z$  exchange.



At NuTeV  $\nu_\mu$  and  $\bar{\nu}_\mu$  beams of a mean energy of 125 GeV were scattered off a target detector and the ratios

$$R^\nu = \frac{\sigma_{\text{NC}}^\nu}{\sigma_{\text{CC}}^\nu},$$

$$R^{\bar{\nu}} = \frac{\sigma_{\text{NC}}^{\bar{\nu}}}{\sigma_{\text{CC}}^{\bar{\nu}}}$$

were measured.

NuTeV measured also the weak mixing angle [NuTeV '02].

$$\sin^2 \theta_w^{\text{on-shell}} = 0.2277 \pm 0.0013 \pm 0.0009$$

→ This is about  $3\sigma$  below the SM prediction !

But the measurement is **indirect**, using the measurements of  $R^\nu, R^{\bar{\nu}}$  and making use of the Paschos-Wolfenstein relation

$$R^- = \frac{R^\nu - rR^{\bar{\nu}}}{1 - r} = \frac{1}{2} - \sin^2 \theta_w + \dots, \quad r = \frac{\sigma_{\text{CC}}^{\bar{\nu}}}{\sigma_{\text{CC}}^\nu} \approx \frac{1}{2}.$$

More precisely, ratios of counting rates

$$R_{\text{exp}}^{\nu} = \frac{\# \text{ of NC-like } \nu \text{ events}}{\# \text{ of CC-like } \nu \text{ events}} \approx R^{\nu}, \quad R_{\text{exp}}^{\bar{\nu}} = \frac{\# \text{ of NC-like } \bar{\nu} \text{ events}}{\# \text{ of CC-like } \bar{\nu} \text{ events}} \approx R^{\bar{\nu}}$$

are measured.

$R_{\text{exp}}^{\nu}, R_{\text{exp}}^{\bar{\nu}}$  can be related to  $R^{\nu}, R^{\bar{\nu}}$  by a detailed MC physics simulation.

The deviation from the SM in terms of  $R_{\text{exp}}^{\nu}, R_{\text{exp}}^{\bar{\nu}}$  are [NuTeV '02]:

$$\Delta R^{\nu} = R_{\text{exp}}^{\nu} - R_{\text{exp}}^{\nu}(SM) = -0.0032 \pm 0.0013 ,$$

$$\Delta R^{\bar{\nu}} = R_{\text{exp}}^{\bar{\nu}} - R_{\text{exp}}^{\bar{\nu}}(SM) = -0.0016 \pm 0.0028 .$$

→  $\Delta R^{\nu, \bar{\nu}}$  : simple starting point for studying MSSM radiative corrections

## – possible explanations

- statistical fluctuation, errors underestimated ?

- re-analyses of EW rad. corr.

[Diener, Dittmaier, Hollik '03 & '05; Arbuzov, Bardin, Kalinovskaya '03]

- relevant SM effects neglected ?

- asymmetry of strange sea-quarks in the nucleon ( $s \neq \bar{s}$ )

- isospin violation ( $u_p \neq d_n$ )

- nuclear effects

- etc. ...

- new physics ?

- modified gauge boson interactions (e.g. in extra dimensions)

- non-renormalizable operators (suppressed by powers of  $\Lambda_{\text{new physics}}^{-1}$ )

- leptoquarks (e.g.  $R$  parity violating SUSY)

- SUSY loop effects (e.g. in MSSM)

- etc. ...

Although the NuTeV "anomaly" is far from being settled, it is interesting, if the MSSM *could* account for such an effect.

Earlier Studies:

- Davidson et al. ['02]
  - rough study in terms of oblique corrections (i.e. momentum transfer  $q = 0$ )
  - no Parton Distribution Functions (PDFs) used
- Kurylov, Ramsey-Musolf, Su ['04] :
  - detailed parameter dependence studied
  - momentum transfer  $q = 0$  approximation
  - no PDFs used

results so far: MSSM not responsible (size ok, but wrong sign)

- our calculation: try to include kinematic effects [OBr, Koch, Hollik]
  - full  $q^2$ -dependence
  - use PDFs
  - use NuTeV cuts on hadronic Energy in final state
  - use mean neutrino beam energy (125 GeV)

- MSSM radiative corrections to  $\nu_\mu N$  DIS

– definition of  $\delta R^{\nu, \bar{\nu}} = R_{\text{MSSM}}^{\nu, \bar{\nu}} - R_{\text{SM}}^{\nu, \bar{\nu}}$

The difference between MSSM and SM prediction,  $\delta R^n = R_{\text{MSSM}}^n - R_{\text{SM}}^n$  with  $R^n = \sigma_{\text{NC}}^n / \sigma_{\text{CC}}^n$  ( $n = \nu, \bar{\nu}$ ), using

$$(\sigma_{\text{NC}}^n)_{\text{NLO}} = (\sigma_{\text{NC}}^n)_{\text{LO}} + \delta\sigma_{\text{NC}}^n \quad (n = \nu, \bar{\nu})$$

$$(\sigma_{\text{CC}}^n)_{\text{NLO}} = (\sigma_{\text{CC}}^n)_{\text{LO}} + \delta\sigma_{\text{CC}}^n \quad (n = \nu, \bar{\nu})$$

can be expanded in  $\delta\sigma_{\text{NC}}^n$  and  $\delta\sigma_{\text{CC}}^n$

$$\delta R^n = \left( \frac{\sigma_{\text{NC}}^n}{\sigma_{\text{CC}}^n} \right)_{\text{LO}} \left( \frac{(\delta\sigma_{\text{NC}}^n)_{\text{MSSM}} - (\delta\sigma_{\text{NC}}^n)_{\text{SM}}}{(\sigma_{\text{NC}}^n)_{\text{LO}}} - \frac{(\delta\sigma_{\text{CC}}^n)_{\text{MSSM}} - (\delta\sigma_{\text{CC}}^n)_{\text{SM}}}{(\sigma_{\text{CC}}^n)_{\text{LO}}} \right).$$

→ Only differences between MSSM and SM radiative corrections and LO cross sections appear in  $\delta R^n$ .



Because of  $R$  parity conservation in the MSSM:

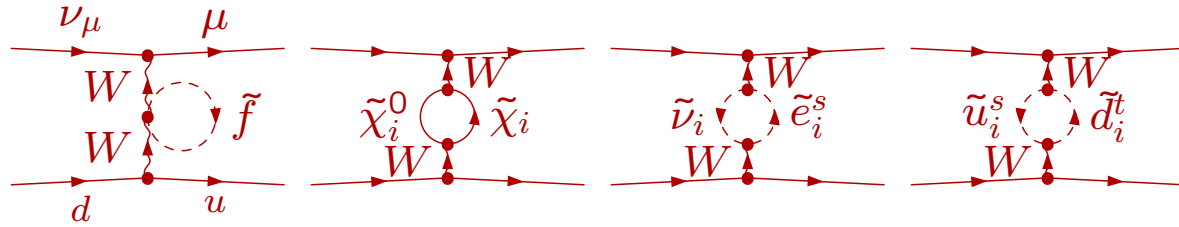
- Born cross section : SM = MSSM (very good approx.)
- real photon emission corrections : SM = MSSM (very good approx.)
- SM = MSSM for SM-like 1-loop graphs without virtual Higgs

Thus:

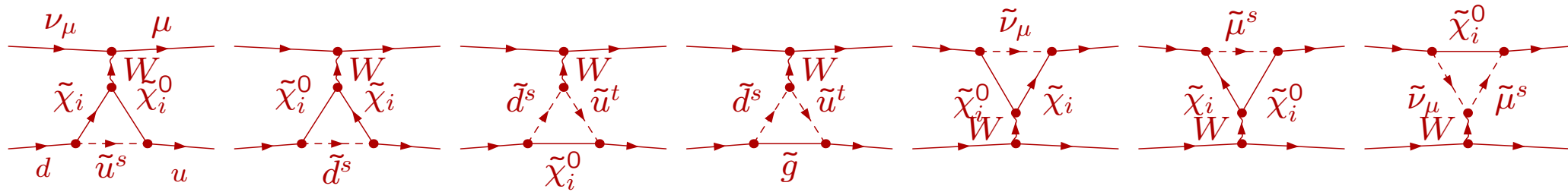
$$\delta\sigma_{\text{MSSM}} - \delta\sigma_{\text{SM}} = \text{const.} \times \left( [\text{superpartner loops}] \right. \\ \left. + [\text{Higgs graphs MSSM} - \text{Higgs graphs SM}] \right) .$$

– superpartner loop corrections

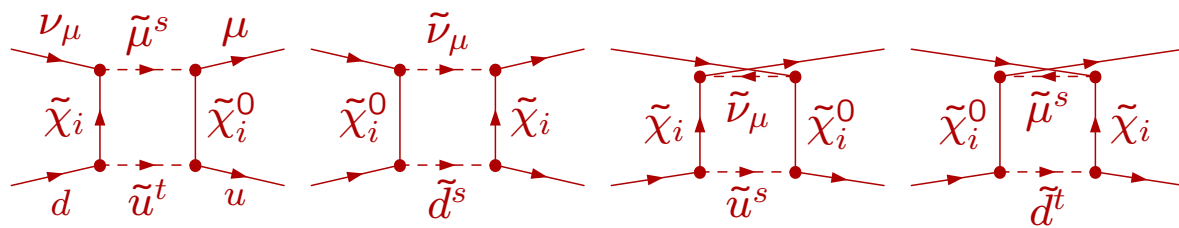
CC self energy insertions



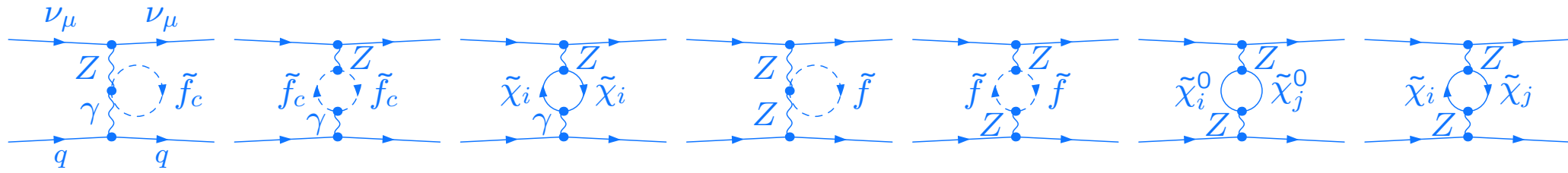
CC vertex corrections



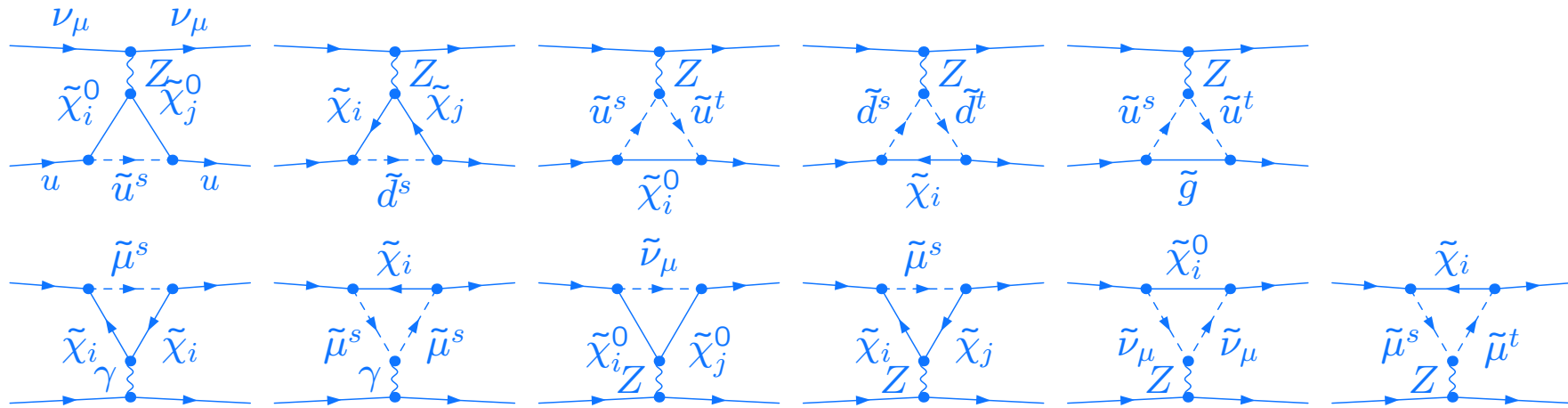
CC box corrections



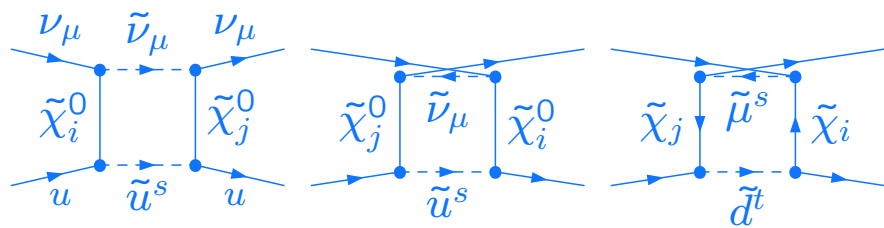
## NC self energy insertions



## NC vertex corrections

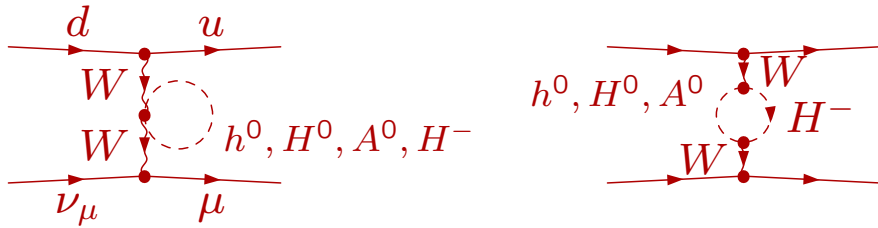


## NC box corrections

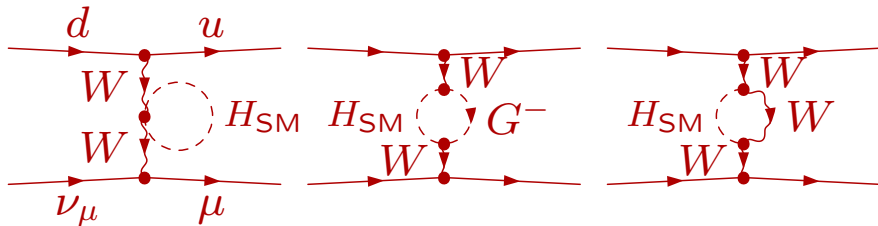


– MSSM-SM Higgs loop difference

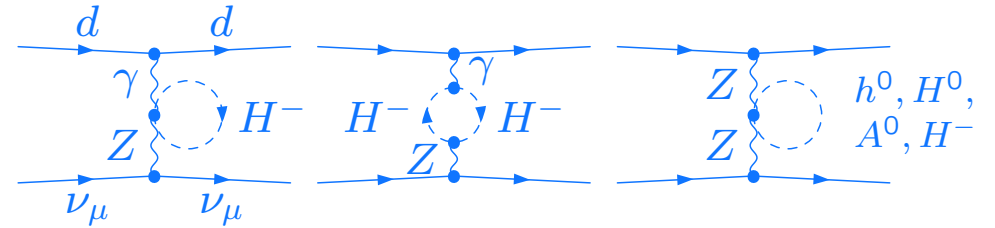
+ CC MSSM Higgs loops



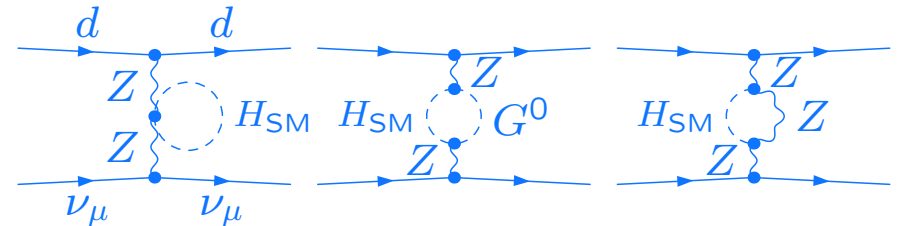
– CC SM Higgs loops



+ NC MSSM Higgs loops



– NC SM Higgs loops



The partonic processes were calculated using  
FeynArts/FormCalc.

[Küblbeck, Böhm, Denner'90],

[Eck '95], [Hahn, Perez-Victoria '99], [Hahn '01],

[Hahn, Schappacher '02]

see : [www.feynarts.de](http://www.feynarts.de)

## • Results for $\delta R^\nu, \delta R^{\bar{\nu}}$

– how to scan over MSSM parameters?

**goal** : find regions of parameter space, where the MSSM might explain the NuTeV anomaly.

→ difficult, large dimensionality of MSSM parameter space

**scanning strategy** : "adaptive scan" [OBr'04]:

→ exploit **adaptive integration by importance sampling**

**method**: calculate an approximation to the integral

$$I = \int_{M_1^{\min}}^{M_1^{\max}} dM_1 \cdots \int_{\tan \beta^{\min}}^{\tan \beta^{\max}} d \tan \beta \ F(\delta R^{\nu(\bar{\nu})}(M_1, \dots, \tan \beta), M_1, \dots, \tan \beta)$$

with VEGAS and **store the sampled parameter points**.

→ automatically, the sample points will be enriched in the regions where  $F(\delta R^{\nu(\bar{\nu})}(M_1, \dots, \tan \beta), M_1, \dots, \tan \beta)$  is large.

some sample choices of  $F$ :

- $F = \begin{cases} 1 & \text{if parameters } (M_1, \dots, \tan \beta) \text{ not excluded} \\ 0 & \text{elsewhere} \end{cases}$

→ enrich points in allowed region

- $F = \begin{cases} \delta R^{\nu(\bar{\nu})}(M_1, \dots, \tan \beta) & \text{if parameters } (M_1, \dots, \tan \beta) \text{ not excluded} \\ 0 & \text{elsewhere} \end{cases}$

→ enrich points where  $|\delta R^{\nu(\bar{\nu})}|$  is large in allowed region

- $F = \begin{cases} \sqrt{(\delta R^\nu(\dots))^2 + (\delta R^{\bar{\nu}}(\dots))^2} & \text{if } \delta R^\nu \text{ and } \delta R^{\bar{\nu}} < 0 \\ 0 & \text{elsewhere} \end{cases}$

→ enrich points where  $\delta R^\nu, \delta R^{\bar{\nu}} < 0$  and  $\sqrt{\dots}$  is large

## – MSSM parameter scan

## restrictions taken into account:

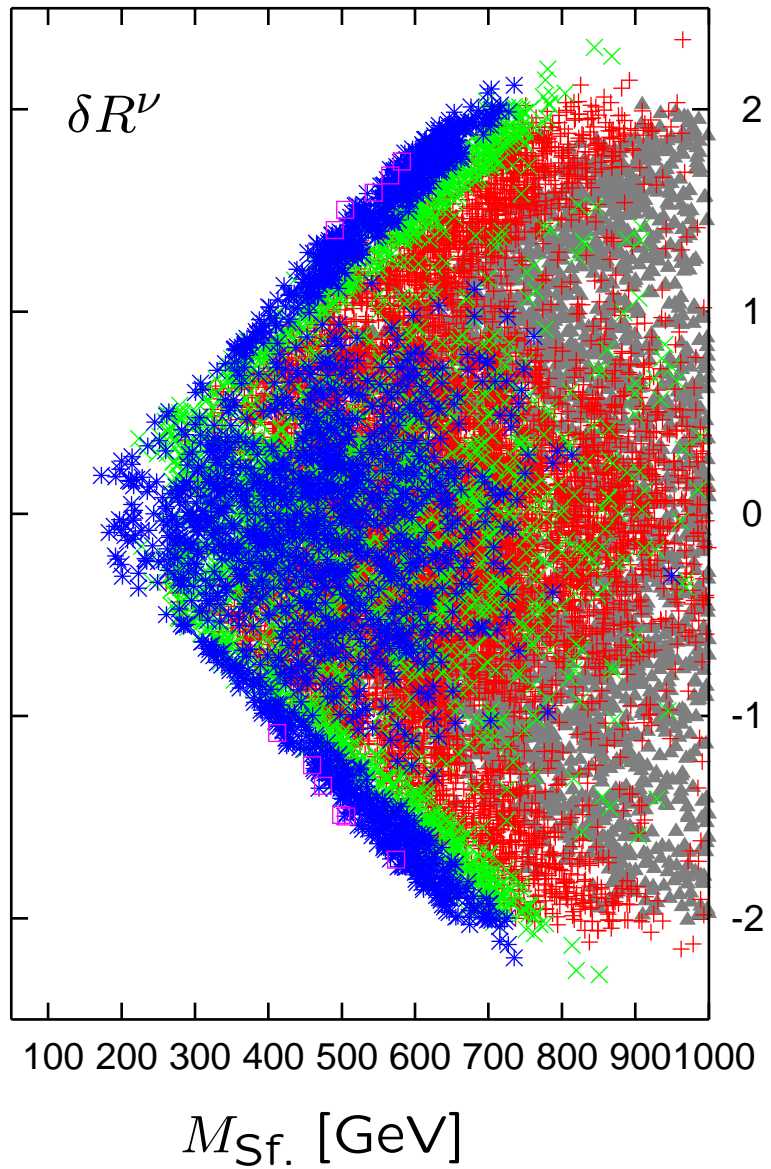
- mass exclusion limits for Higgs bosons and superpartners
- $\Delta\rho$ -constraint on sfermion mixing

## quantities varied:

$M_1, M_2, M_{\text{gluino}}$	:	10 ... 1000 GeV	$\mu$	:	-2000 ... 2000 GeV
$M_{\text{Sf.}}$	:	10 ... 1000 GeV	$A_b, A_t, A_\tau$	:	-2000 ... 2000 GeV
$m_{A^0}$	:	10 ... 1000 GeV	$\tan\beta$	:	1 ... 50



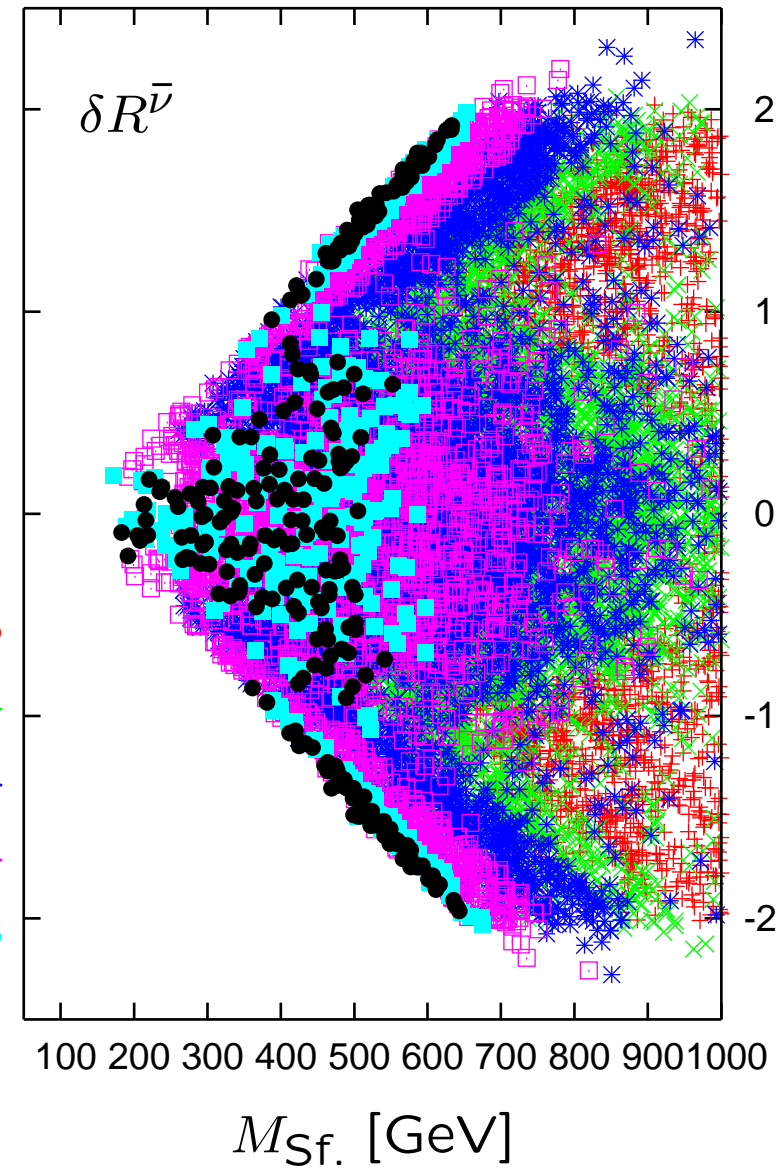
Scan for large values of  $|\delta R^\nu|$  and  $|\delta R^{\bar{\nu}}|$  with parameter restrictions

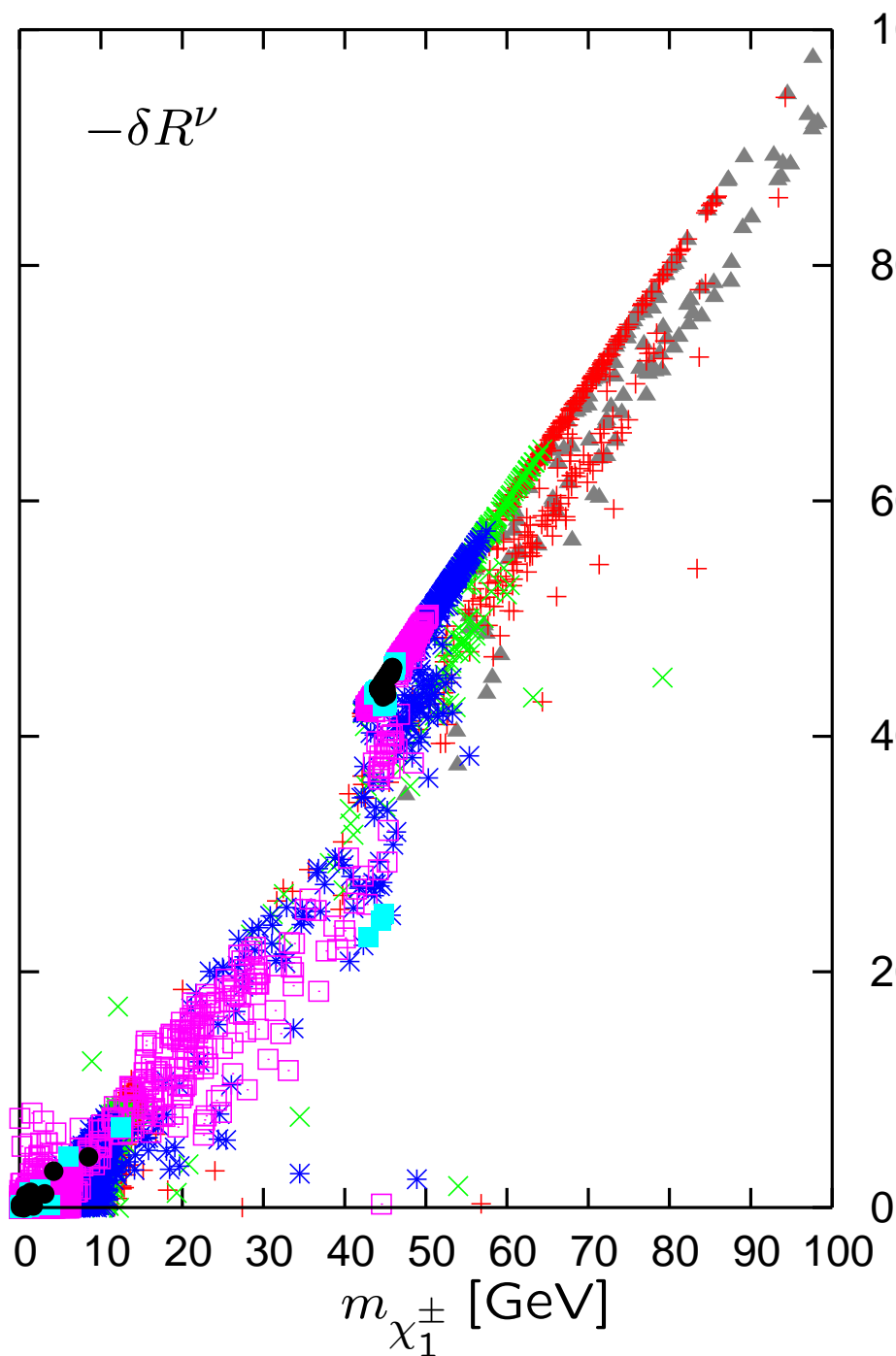


←  $X_t [\text{TeV}]$  →

$\delta R^{\nu(\bar{\nu})}$

- $0 - 1 \cdot 10^{-5}$
- $1 \cdot 10^{-5} - 5 \cdot 10^{-5}$
- $5 \cdot 10^{-5} - 1 \cdot 10^{-4}$
- $1 \cdot 10^{-4} - 3 \cdot 10^{-4}$
- $3 \cdot 10^{-4} - 8 \cdot 10^{-4}$
- $8 \cdot 10^{-4} - 1 \cdot 10^{-3}$
- $> 1 \cdot 10^{-3}$

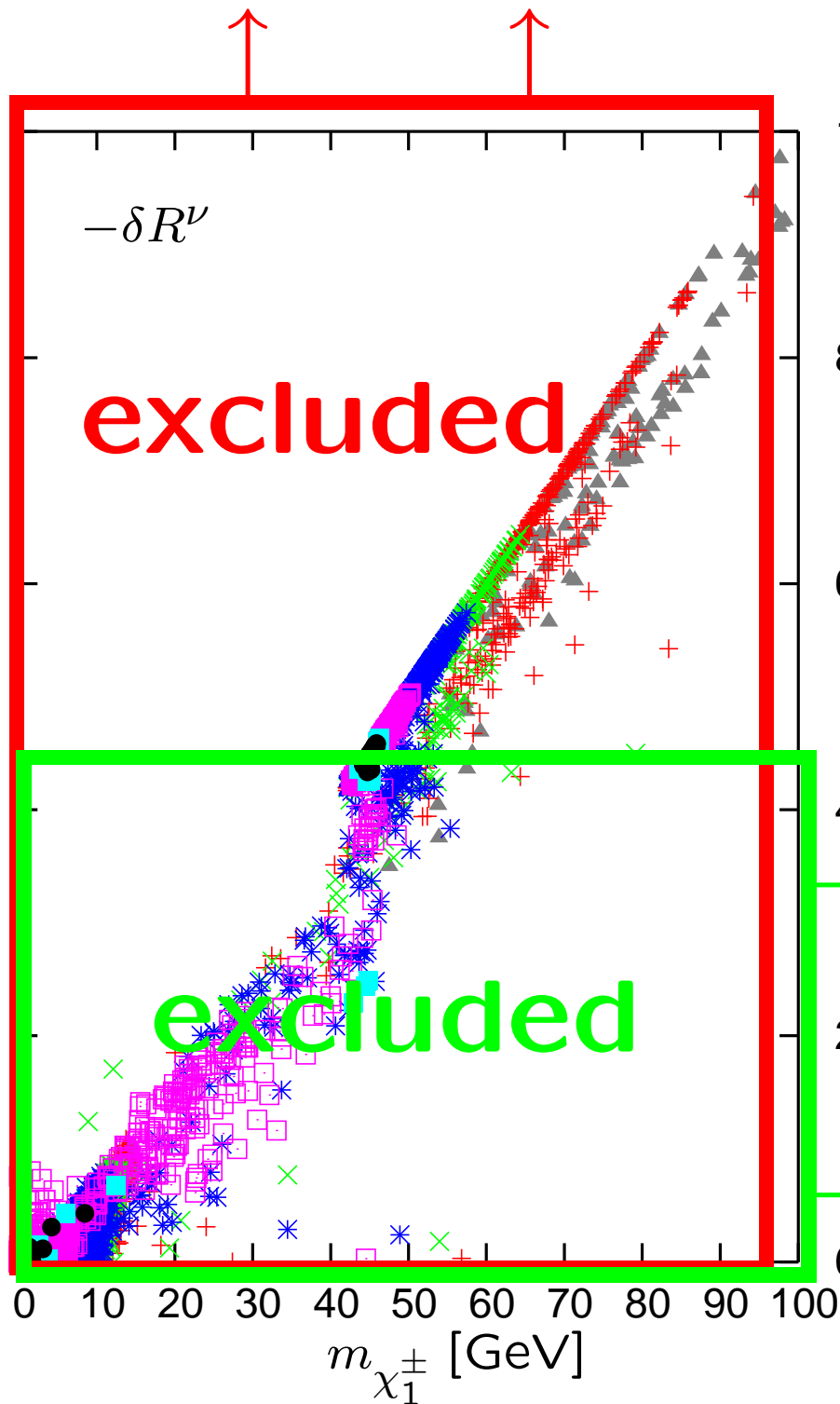




Scan for negative values of  $\delta R^\nu$   
without parameter restrictions

$-\delta R^\nu$

- $0 - 1 \cdot 10^{-5}$
- $1 \cdot 10^{-5} - 5 \cdot 10^{-5}$
- $5 \cdot 10^{-5} - 1 \cdot 10^{-4}$
- $1 \cdot 10^{-4} - 3 \cdot 10^{-4}$
- $3 \cdot 10^{-4} - 8 \cdot 10^{-4}$
- $8 \cdot 10^{-4} - 1 \cdot 10^{-3}$
- $> 1 \cdot 10^{-3}$



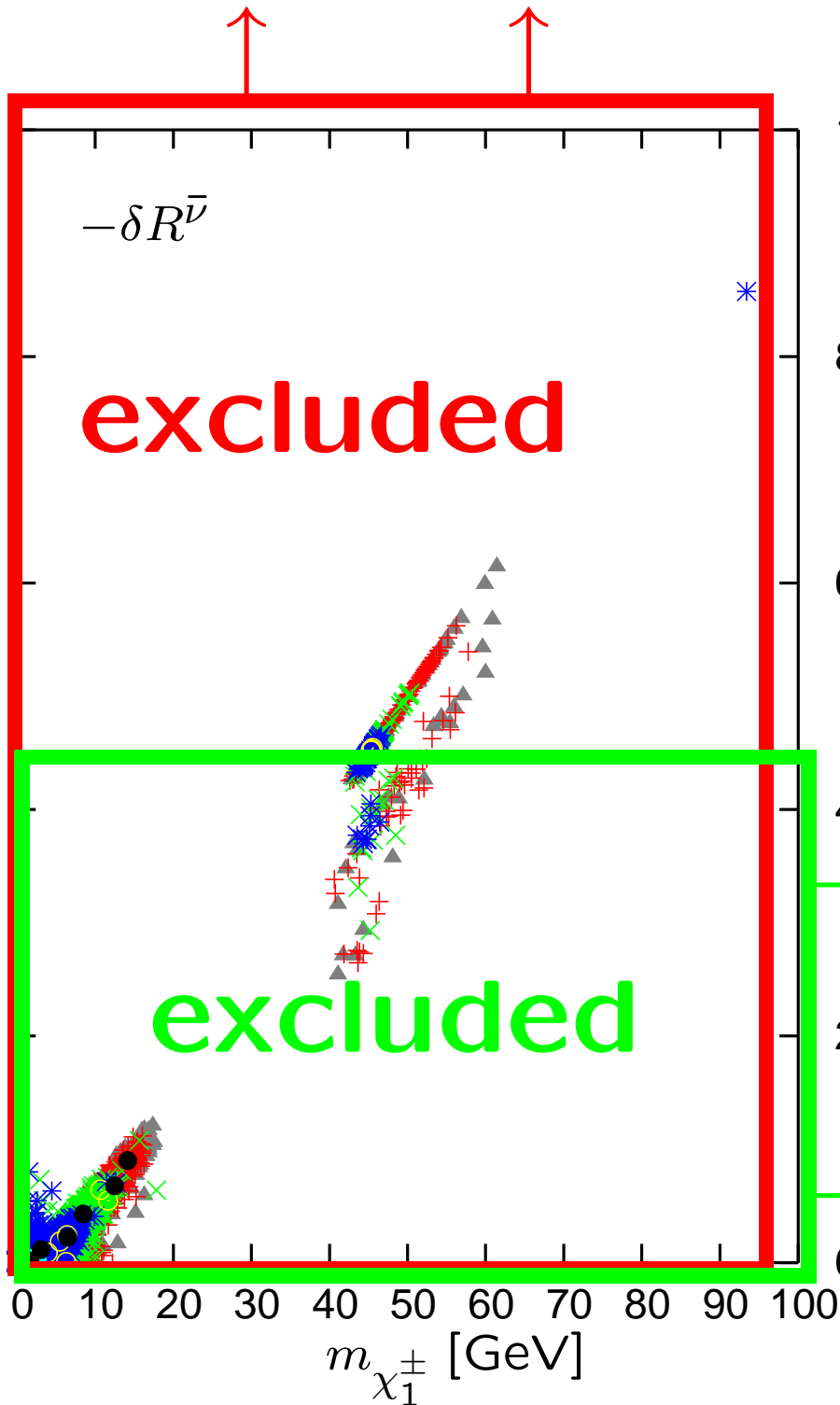
Scan for negative values of  $\delta R^\nu$  without parameter restrictions

$-\delta R^\nu$   
 $0 - 1 \cdot 10^{-5}$   
 $1 \cdot 10^{-5} - 5 \cdot 10^{-5}$   
 $5 \cdot 10^{-5} - 1 \cdot 10^{-4}$   
 $1 \cdot 10^{-4} - 3 \cdot 10^{-4}$   
 $3 \cdot 10^{-4} - 8 \cdot 10^{-4}$   
 $8 \cdot 10^{-4} - 1 \cdot 10^{-3}$   
 $> 1 \cdot 10^{-3}$

$$m_{\chi_1^\pm} > 94 \text{ GeV}$$

$$m_{\chi_1^0} > 46 \text{ GeV}$$

→ all interesting regions excluded



Scan for negative values of  $\delta R^{\bar{\nu}}$  without parameter restrictions

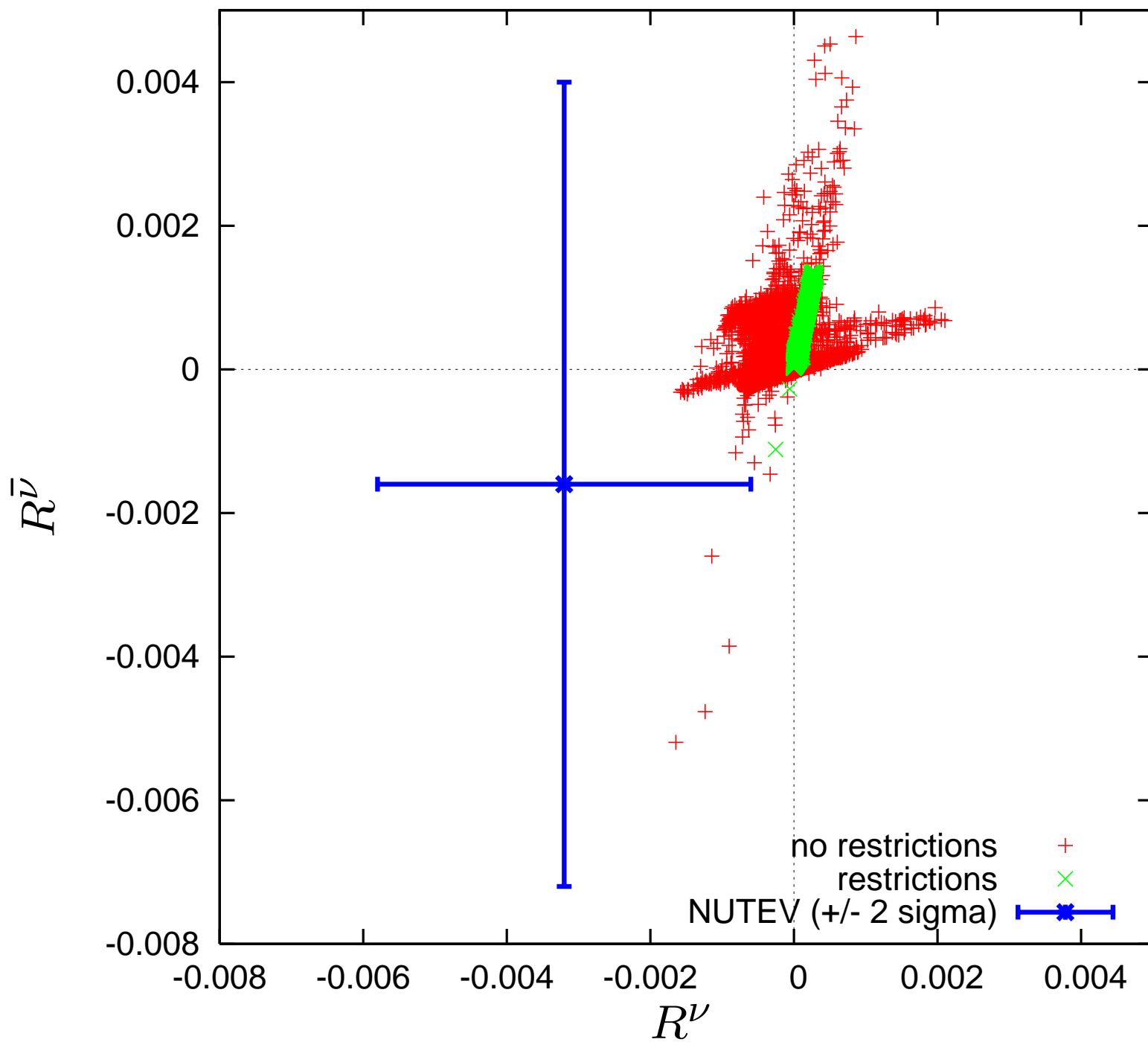
$-\delta R^{\bar{\nu}}$

$0 - 1 \cdot 10^{-5}$   
 $1 \cdot 10^{-5} - 5 \cdot 10^{-5}$   
 $5 \cdot 10^{-5} - 1 \cdot 10^{-4}$   
 $1 \cdot 10^{-4} - 3 \cdot 10^{-4}$   
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 $> 1 \cdot 10^{-3}$

$$m_{\chi_1^\pm} > 94 \text{ GeV}$$

$$m_{\chi_1^0} > 46 \text{ GeV}$$

→ all interesting regions excluded



## Summary

- The NuTeV measurement of  $\sin^2 \theta_w$  is intriguing but has to be further established (especially confirmation by other experiment(s) is desirable).
- Loop effects in the MSSM are not capable of explaining the NuTeV "anomaly". (size can be right, but sign is wrong).
- If the "anomaly" was established, the MSSM would be in trouble.

Our result can be easily combined with the one-loop SM result.  
→ The complete MSSM one-loop prediction for  $\nu N$  scattering is available for future analyses.