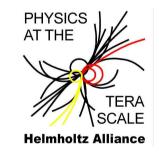
Higgs Bosons and Other New Phenomena

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outline :

- Electroweak Interaction, Higgs Bosons & beyond
 - The Origin of the Electroweak Interaction
 - The Higgs Boson: What is it good for?
 - How to find Higgs Bosons?
 - What else to expect at the LHC?
- Selected Projects
 - SM Higgsstrahlung (NNLO QCD)
 - Randall-Sundrum scalar sector constrained
 - New physics in $\gamma\gamma/WW/ZZ$ production

[Electroweak Interaction, Higgs Bosons & beyond]

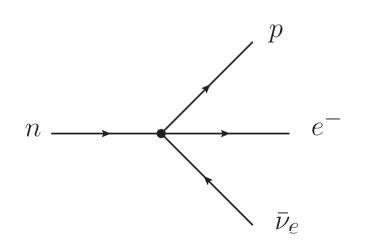
- The Origin of the Electroweak Interaction Beta decay 1911: Hahn, Meitner: observation : $n \rightarrow p \ e^- + missing \ energy$

Puzzle:

- continuous energy spectrum of electrons observed
- discrete spectrum expected (energy difference between n and p state)

Bohr: energy is *really* missing Pauli (1930): $n \rightarrow p e^- +$ neutrino (very weakly interacting)

Fermi (1934): "Fermi Model"



- short-range interaction
- good description for energies well below $G_F^{-1/2} \approx 300 \text{ GeV}$ or equivalently length scales well above $\approx 10^{-18}m$ [= 0.001 × size of atomic nuclei].
- but: bad high energy behaviour

[Electroweak Interaction, Higgs Bosons & beyond]

- The Origin of the Electroweak Interaction Beta decay 1911: Hahn, Meitner: observation : $n \rightarrow p \ e^- + missing \ energy$

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Fermi (1934): "Fermi Model": improvements

Lee, Yang, Wu (1957): Parity violation in weak interactions Marshak, Sudarshan (1957) [Feynman, Gell-Mann]: V - A theory

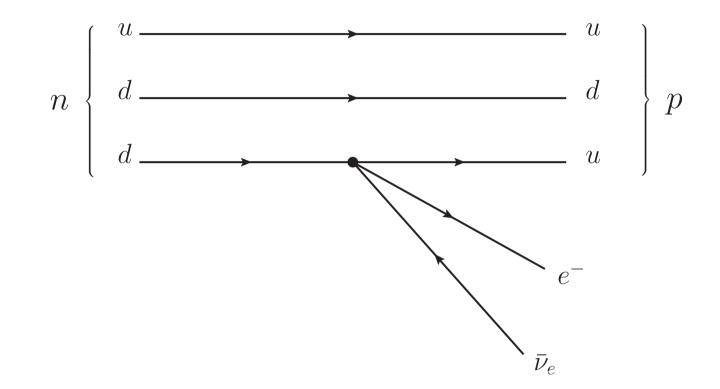
$$\mathcal{L} \propto G_F[\bar{\psi}_A \underbrace{(\gamma_\mu - \gamma_\mu \gamma_5)}_{V-A} \psi_B][\bar{\psi}_C \underbrace{(\gamma^\mu - \gamma^\mu \gamma_5)}_{V-A} \psi_D]$$
$$\propto G_F[\bar{\psi}_A \gamma_\mu P_L \psi_B][\bar{\psi}_C \gamma^\mu P_L \psi_D] \text{ with } P_L = \frac{1}{2}(1 - \gamma_5)$$

• short-range interaction of left-chiral components

[Electroweak Interaction, Higgs Bosons & beyond, Origin of Electroweak Interaction]

Beta decay: current understanding:

Quark parton model [Bjorken, Paschos; Feynman 1969]

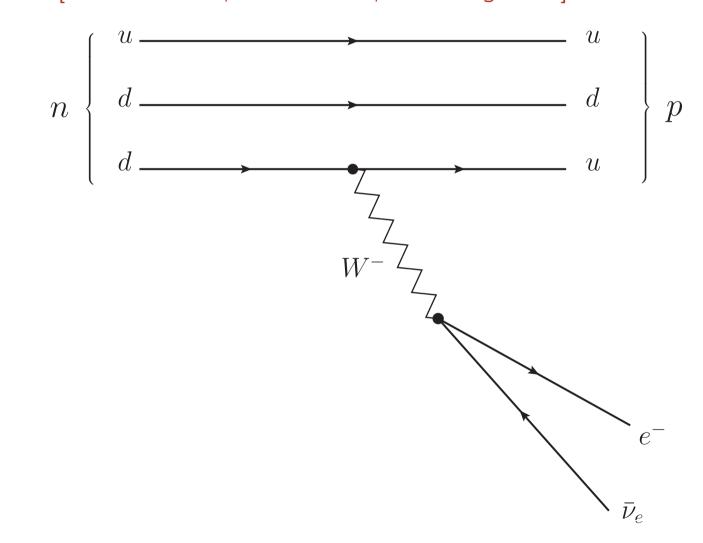


[Electroweak Interaction, Higgs Bosons & beyond, Origin of Electroweak Interaction]

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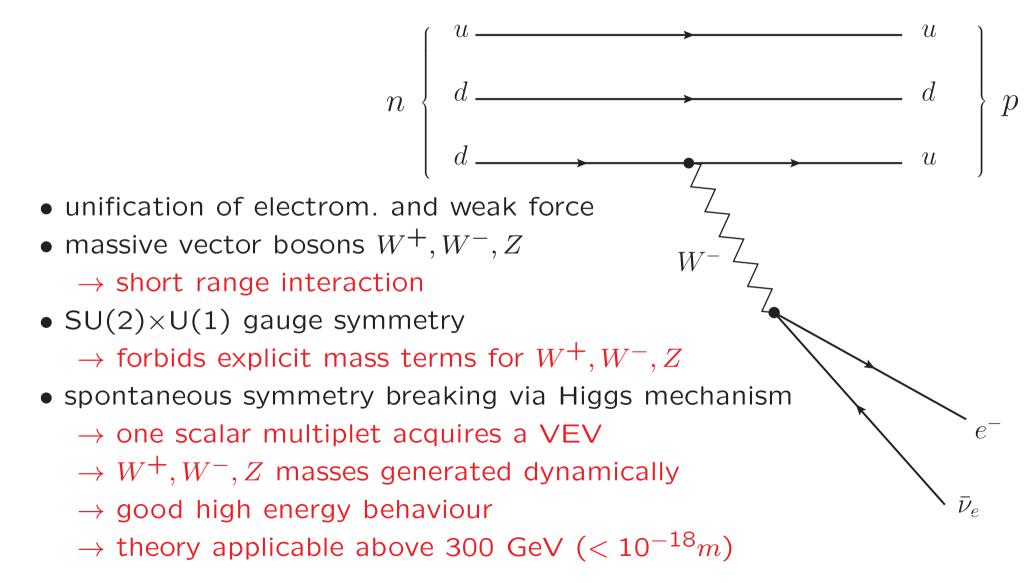
with electroweak interaction [Glashow 1961, Salam 1968, Weinberg 1967]



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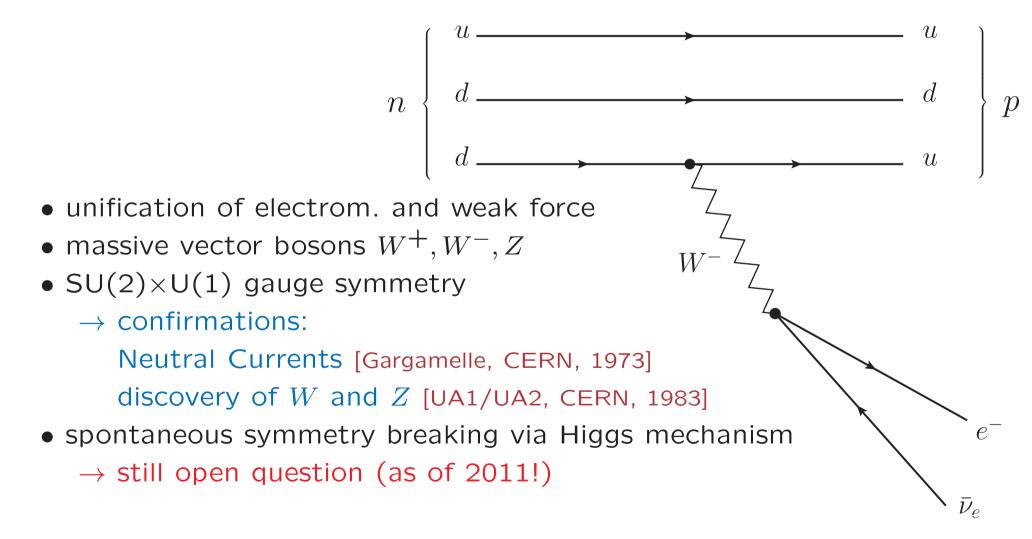
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Beta decay: current understanding:

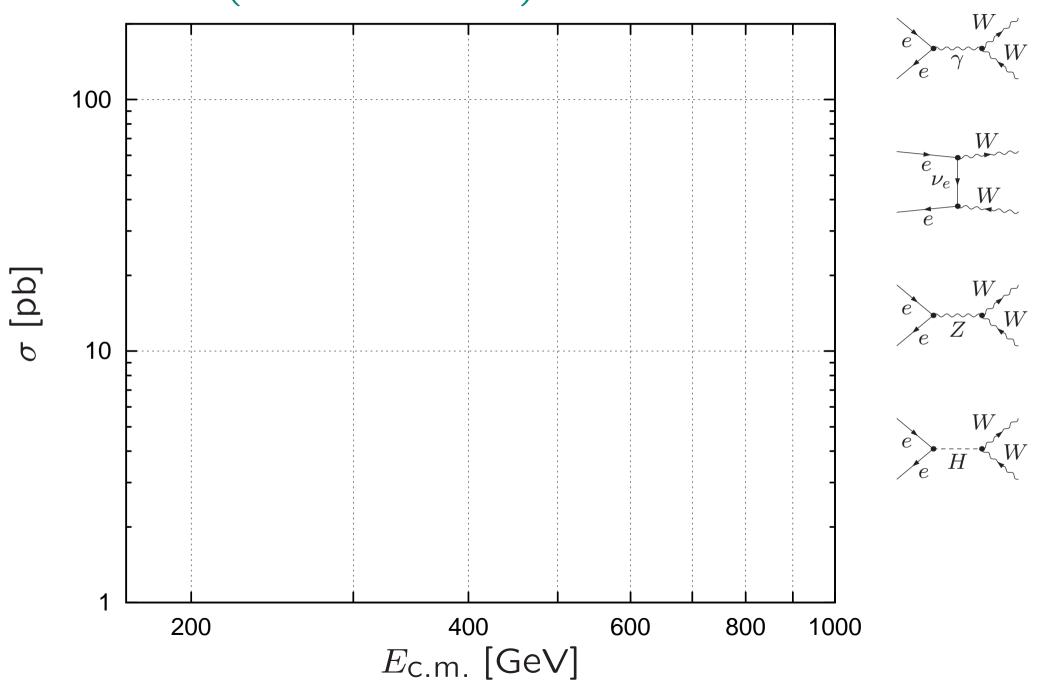
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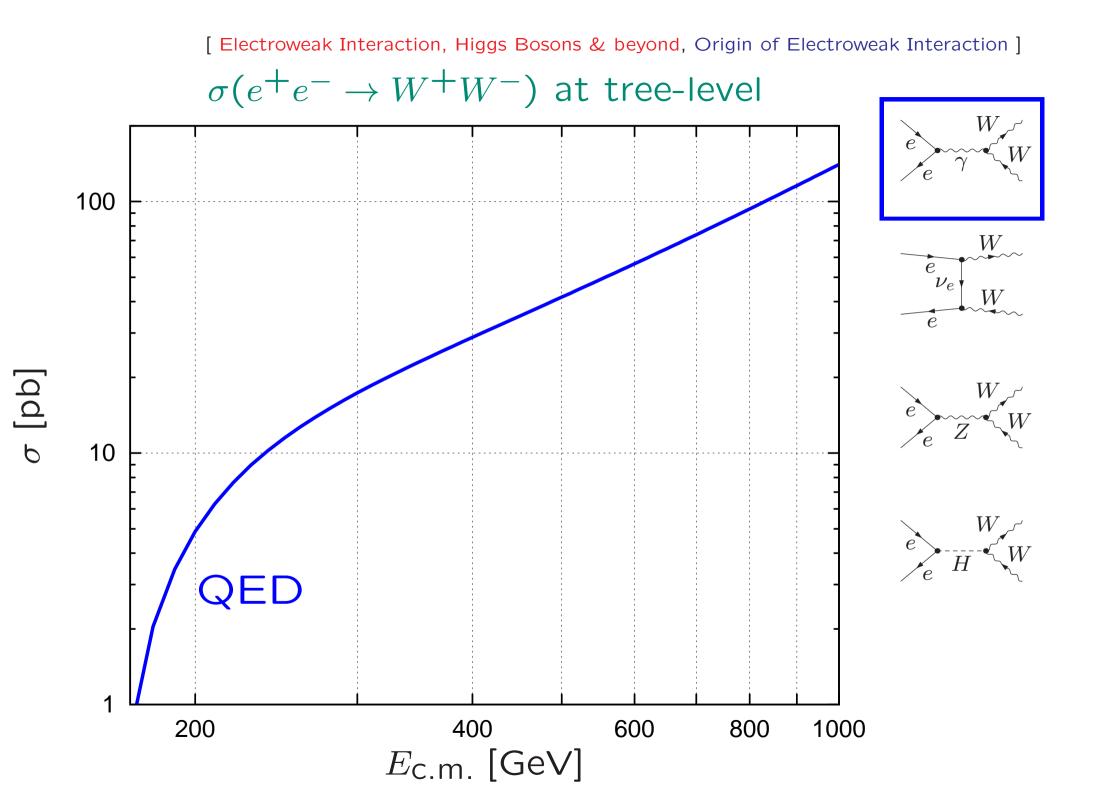


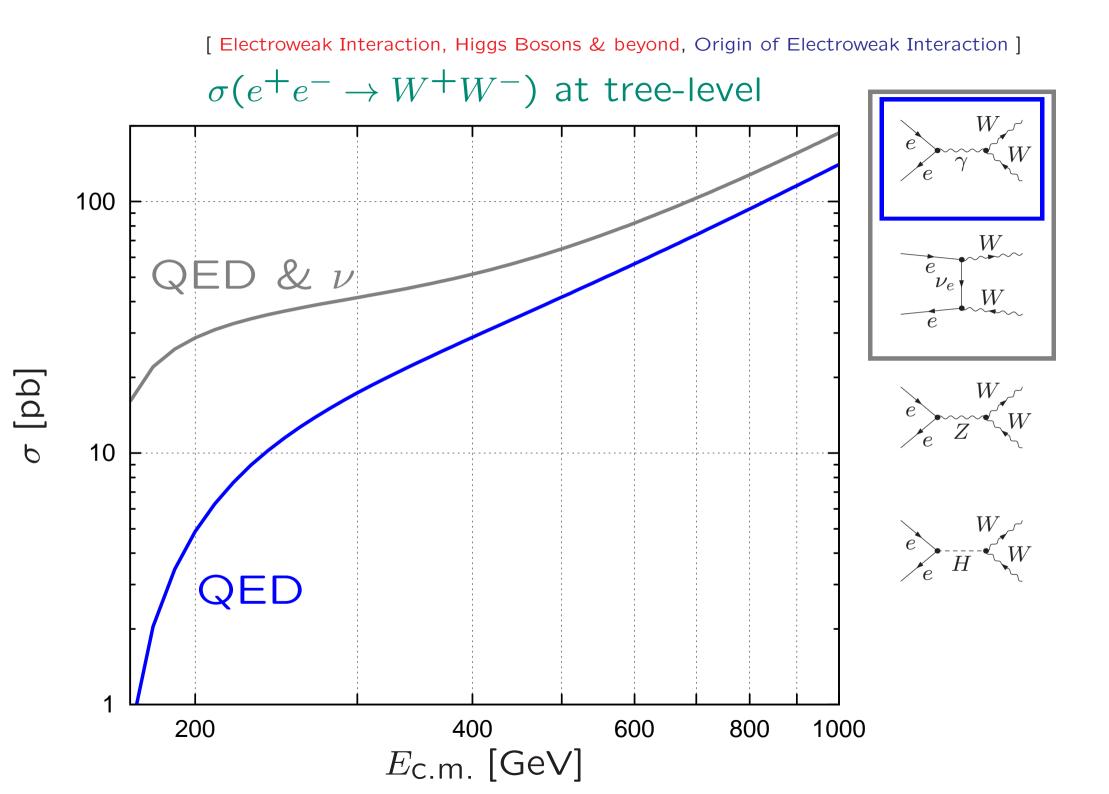
[Electroweak Interaction, Higgs Bosons & beyond, Origin of Electroweak Interaction]

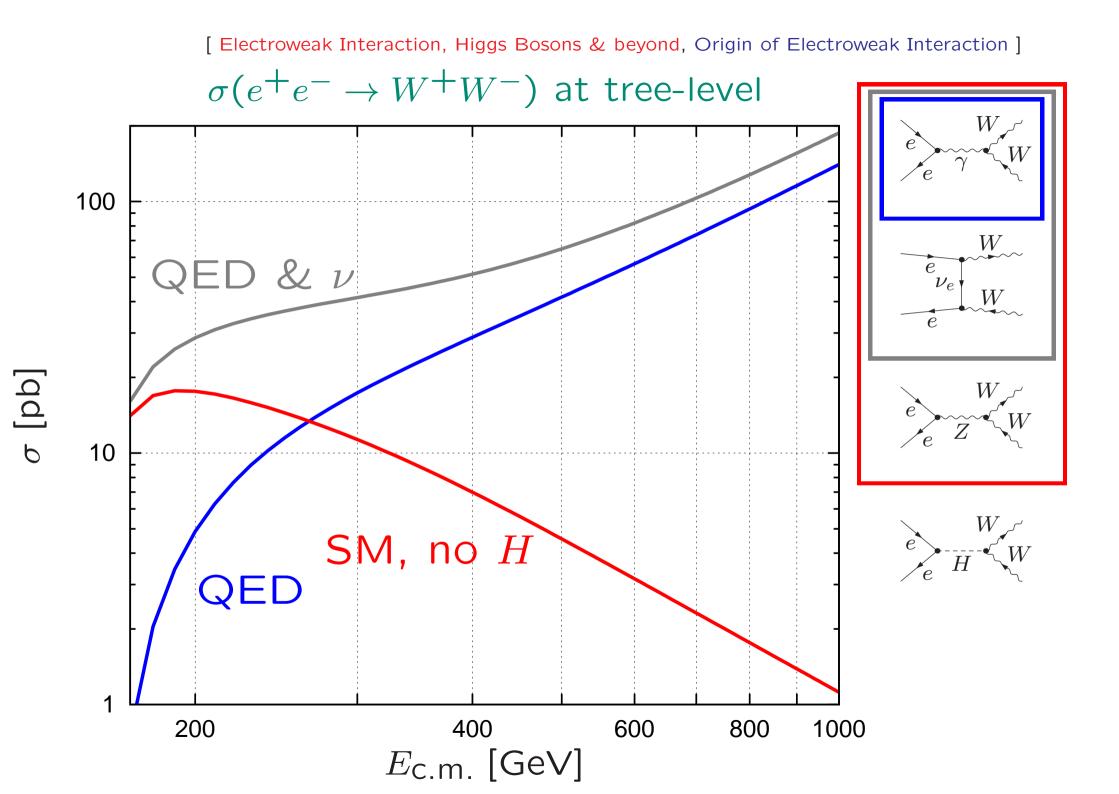
 $\sigma(e^+e^- \rightarrow W^+W^-)$ at tree-level

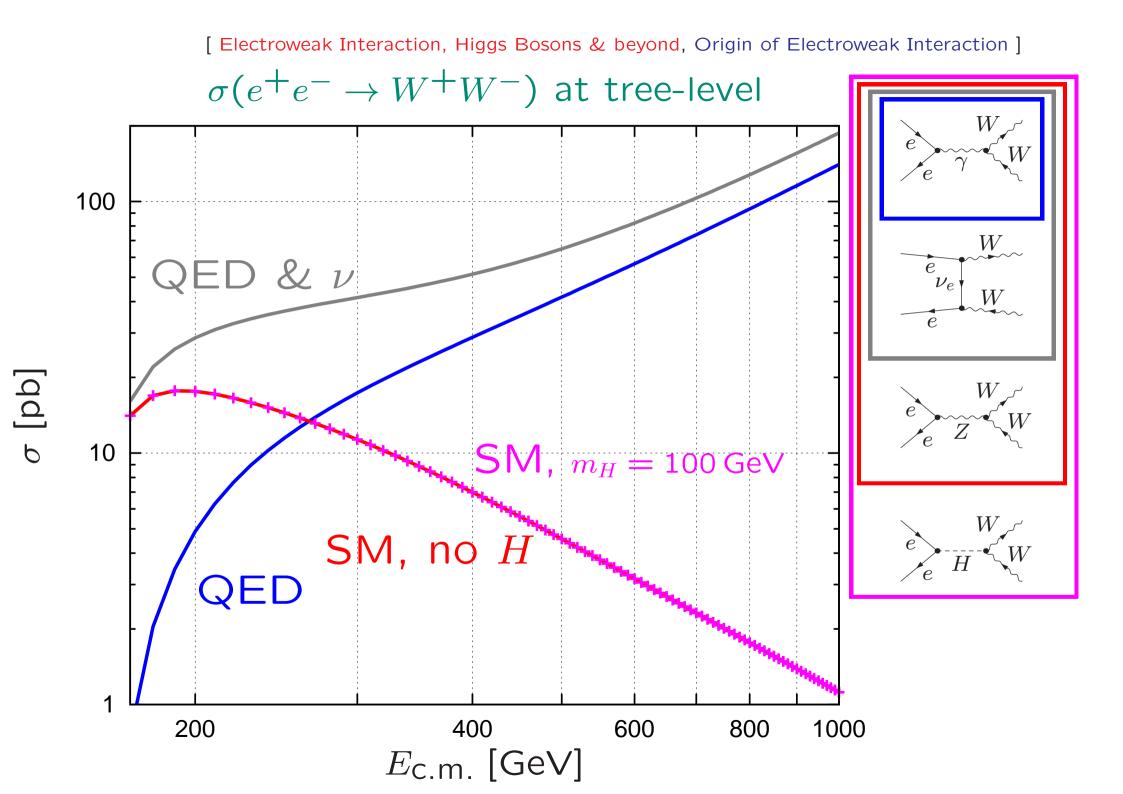


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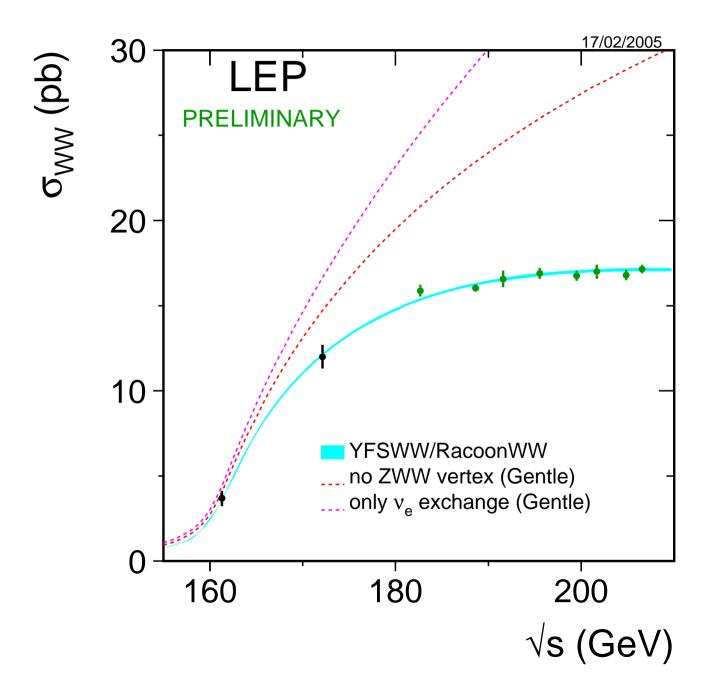






[Electroweak Interaction, Higgs Bosons & beyond, Origin of Electroweak Interaction]

measurement of $\sigma(e^+e^- \rightarrow W^+W^-)$ at LEP 2:



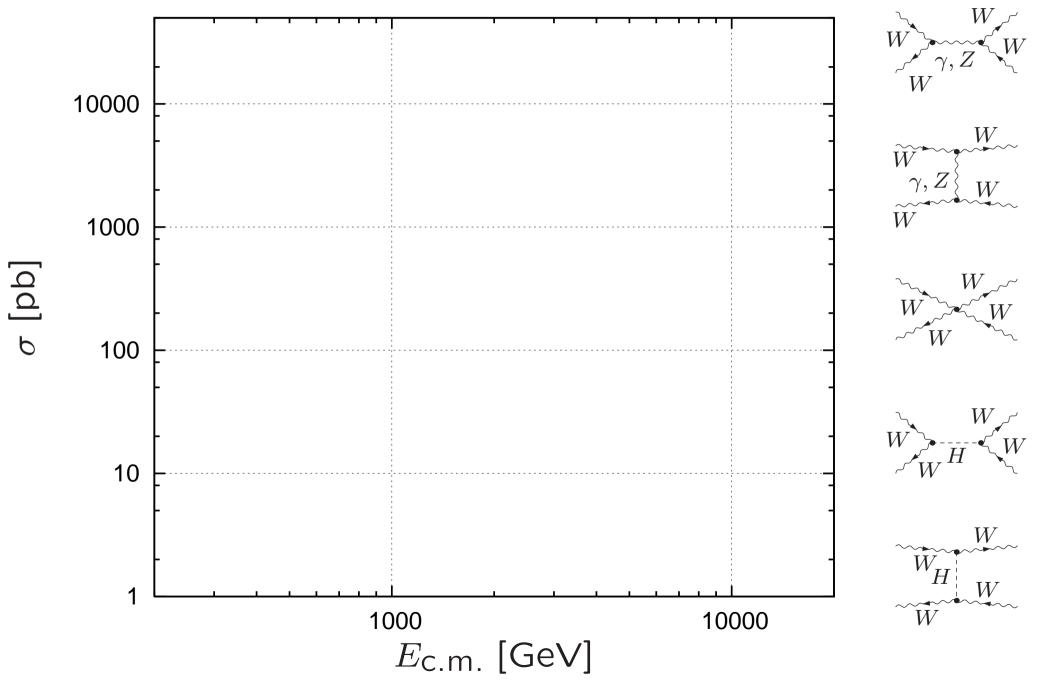
[Electroweak Interaction, Higgs Bosons & beyond]

- The Higgs Boson: What is it good for?
 - The Higgs mechanism (in the electroweak Standard Model):
 - The Higgs field has 4 components & doesn't vanish in the ground state
 - The ground state configuration acts as a medium (background field) with which all particles interact (coupling strength \propto mass)
 - 3 components promote Z, W^+, W^- to massive (3 component) vector particles from massless (2 component) ones
 - 1 component is an additional physical d.o.f. $H \rightarrow$ the Higgs boson (coupling strength to other particles \propto mass)
 - The Higgs gives mass to all elementary particles: (e.g. e^- , q, Z, W^{\pm})
 - the Higgs mechanism is a general concept (choice of Higgs field not unique)
 - it explains *how* masses arise but not *what* mass values

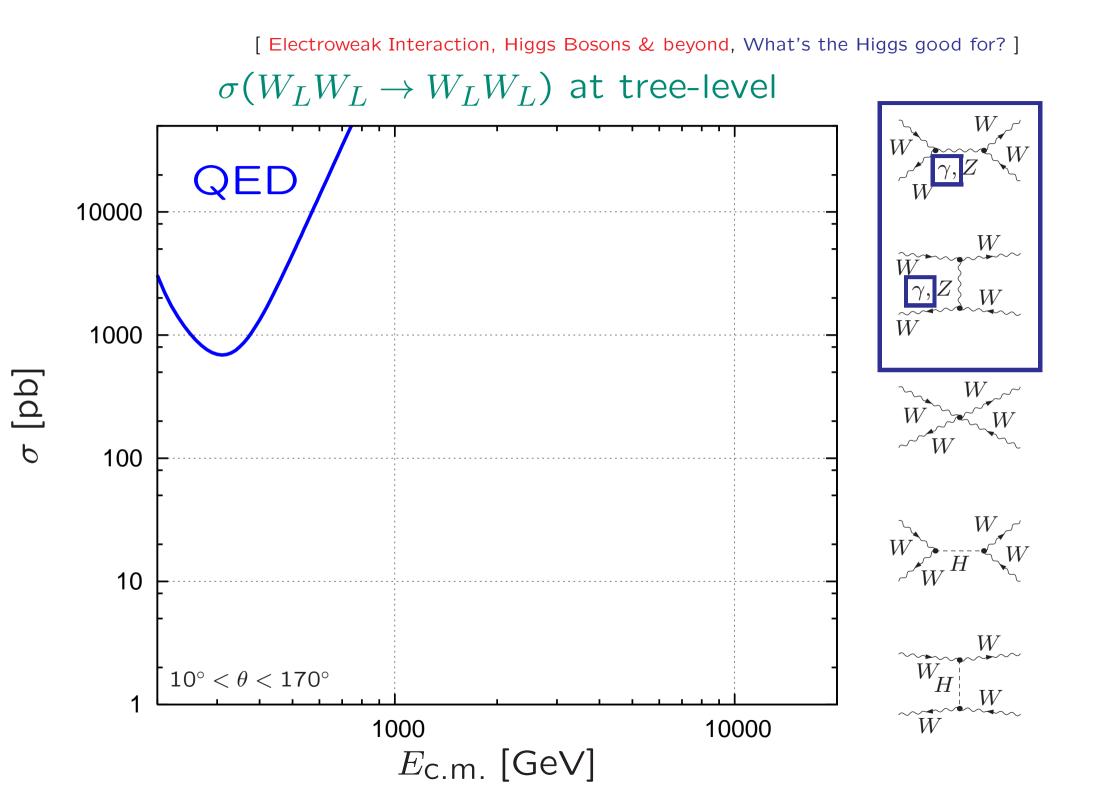
I The Higgs cures bad high energy behaviour: (example $W_L W_L$ scattering)

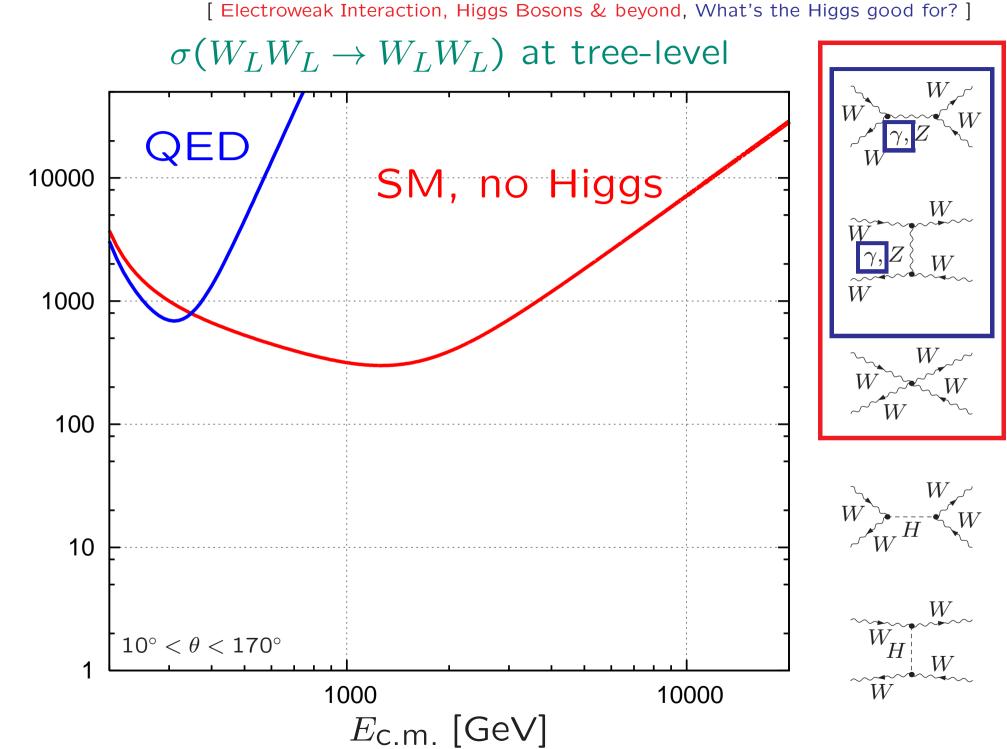
[Electroweak Interaction, Higgs Bosons & beyond, What's the Higgs good for?]

 $\sigma(W_L W_L \to W_L W_L)$ at tree-level

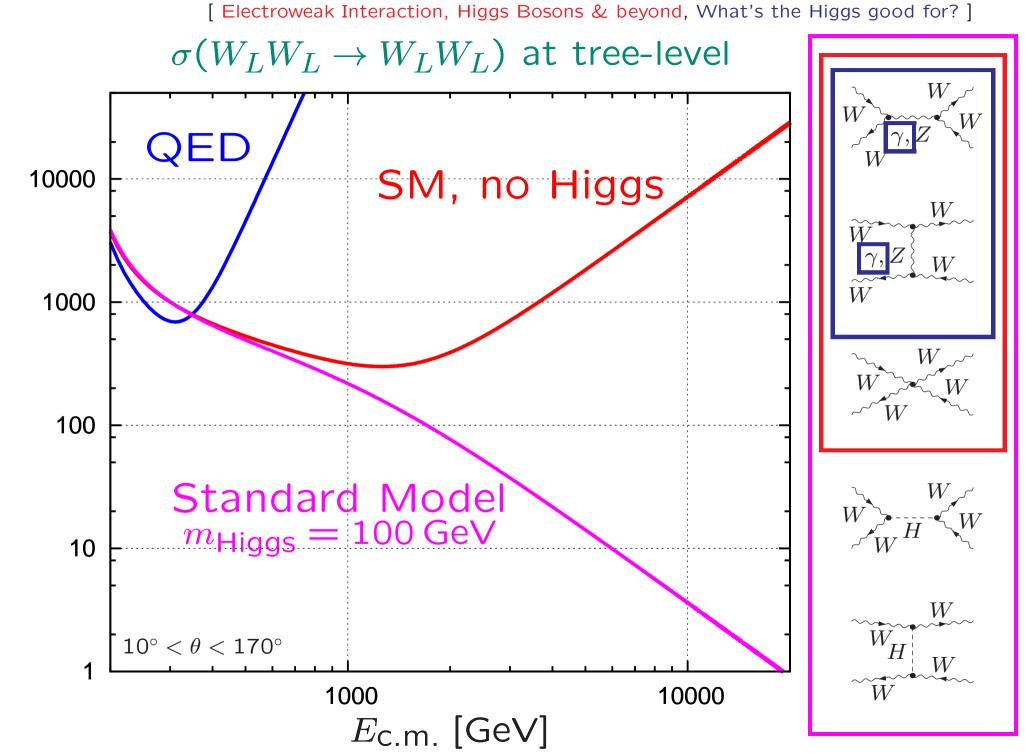


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 σ [pb]

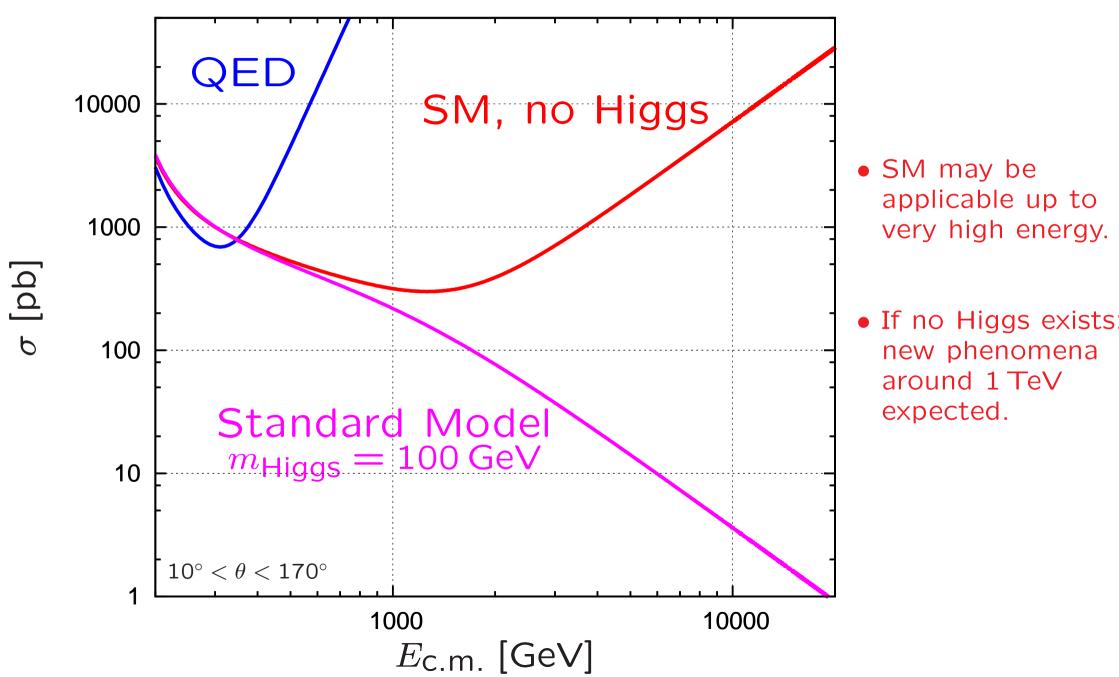


σ [pb]

[Electroweak Interaction, Higgs Bosons & beyond, What's the Higgs good for?]

 $\sigma(W_L W_L \to W_L W_L)$ at tree-level

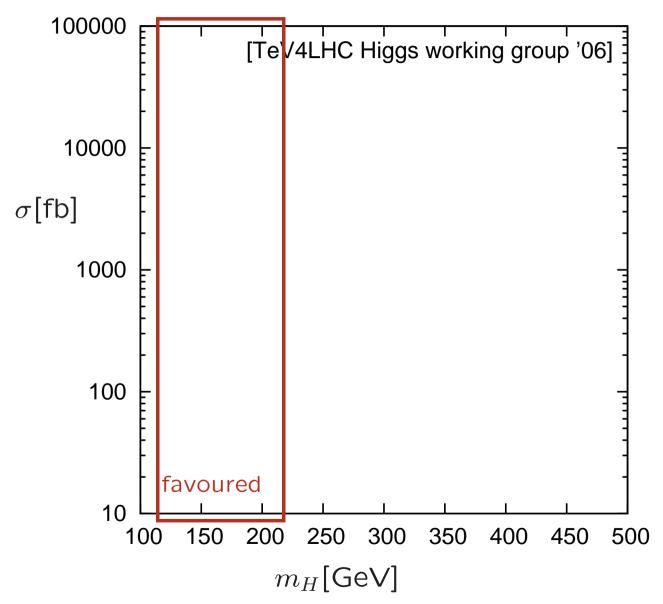
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■ SM Higgs production @ LHC :

 \rightarrow consider:

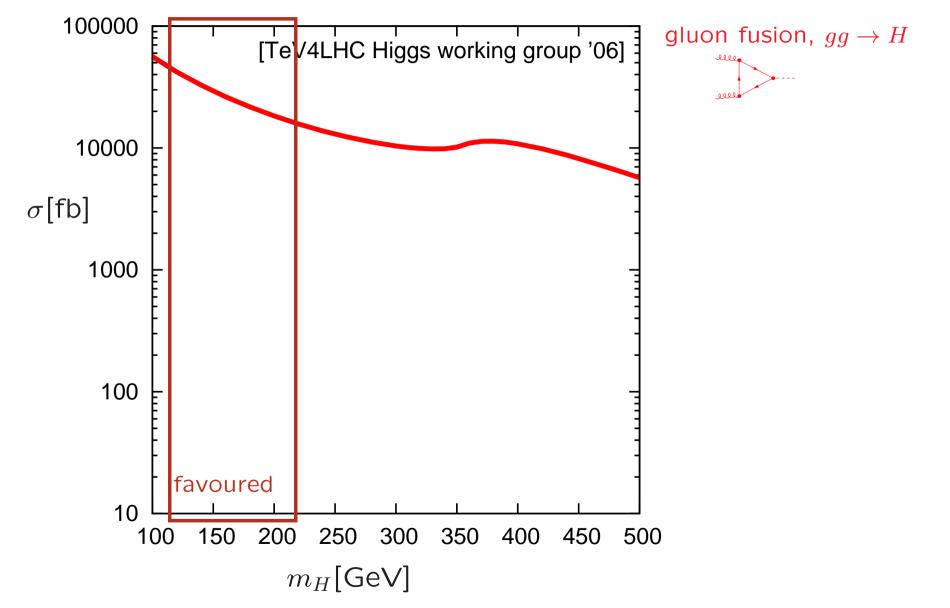
a) Higgs couplings \propto mass. b) Ordinary matter is very light. c) Huge # of gluon collisions.



SM Higgs production @ LHC :

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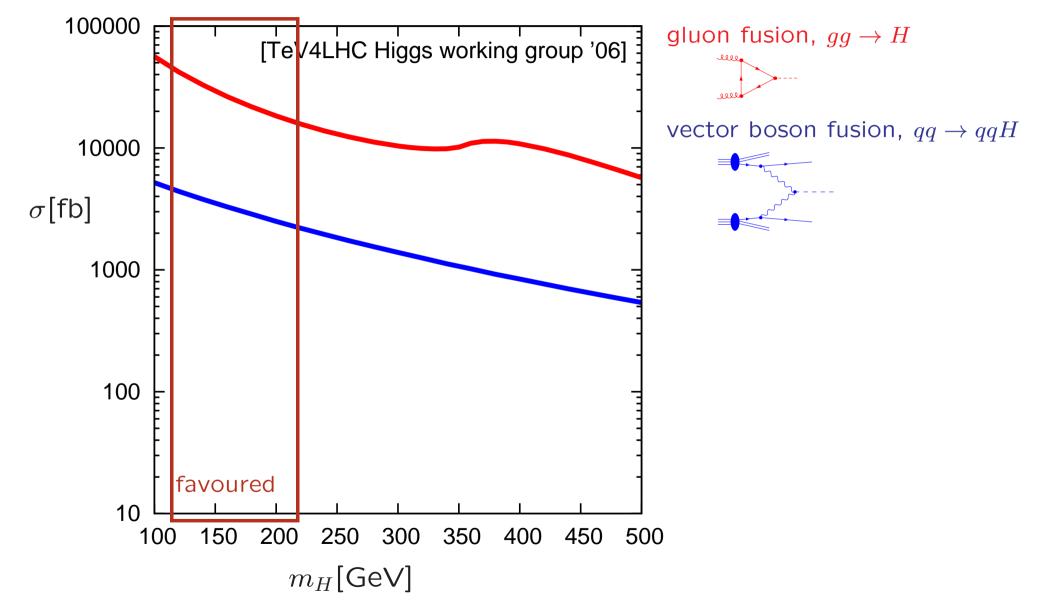
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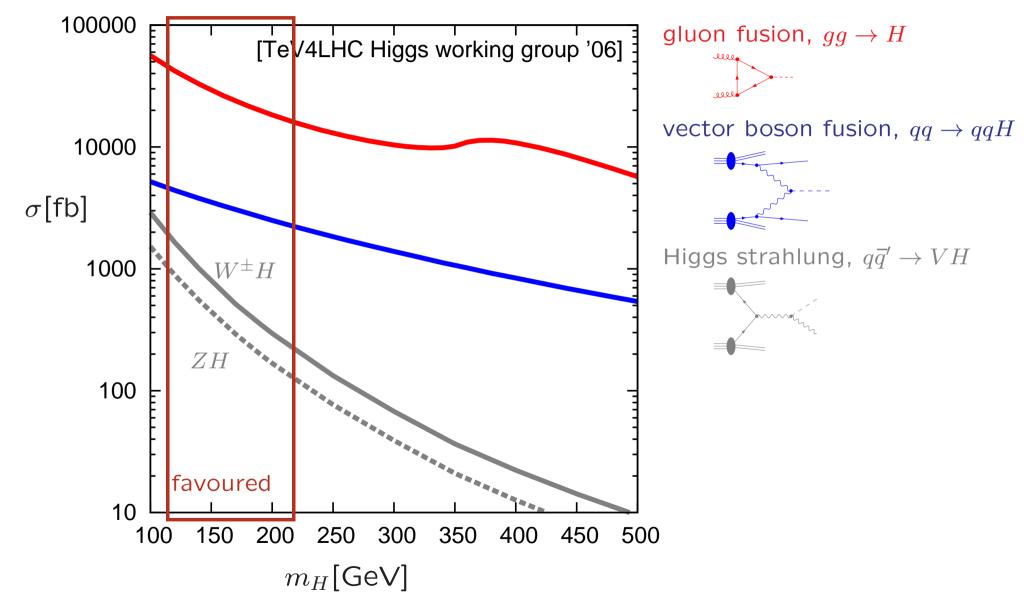
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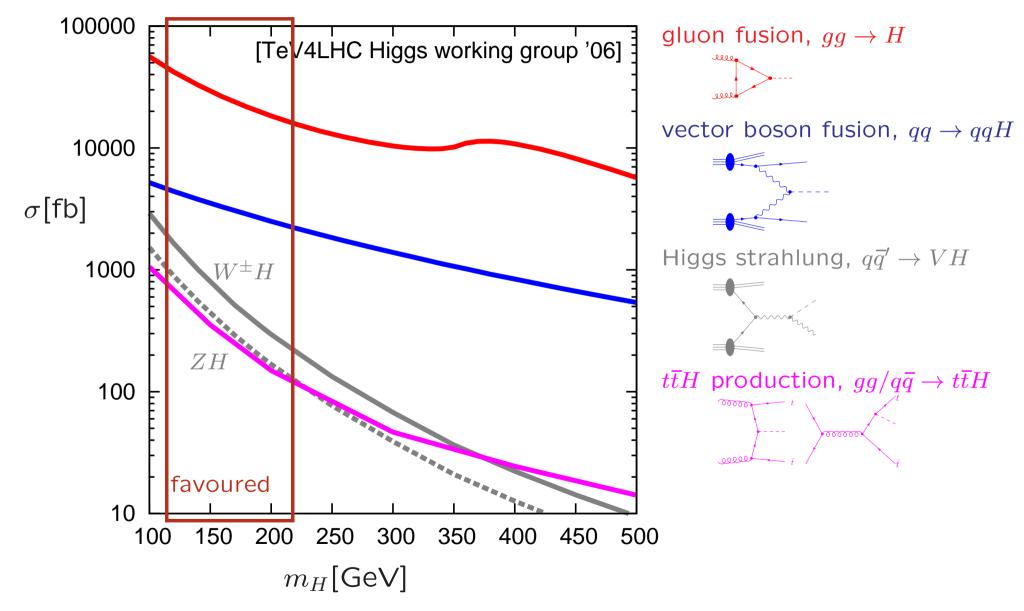
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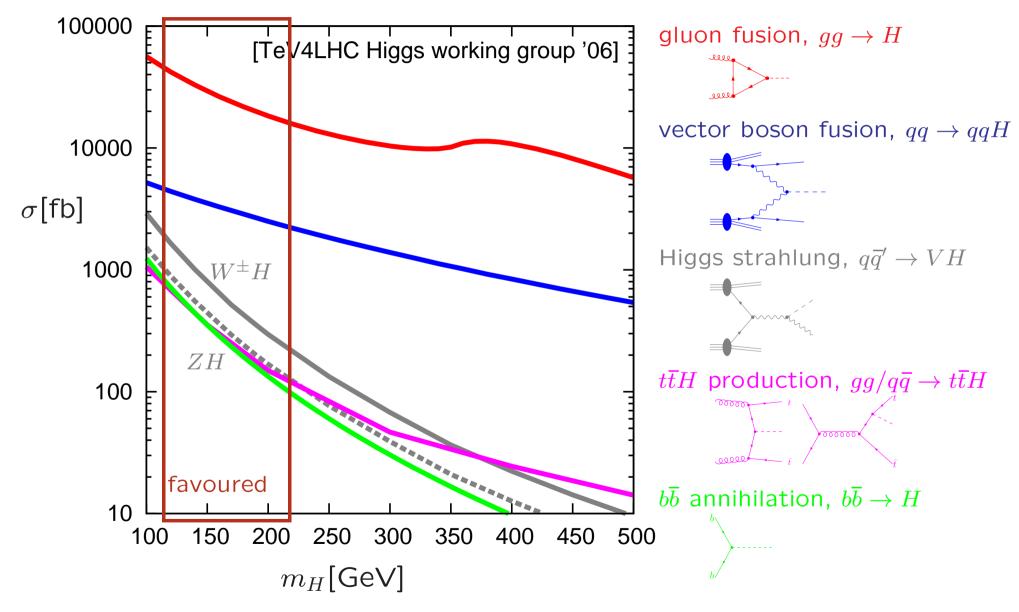
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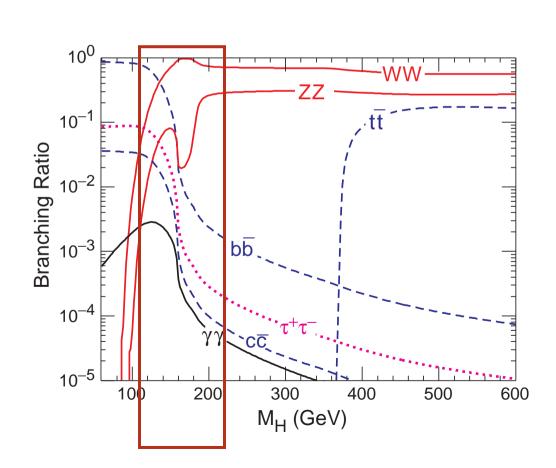
[Electroweak Interaction, Higgs Bosons & beyond, How to find Higgs Bosons?]

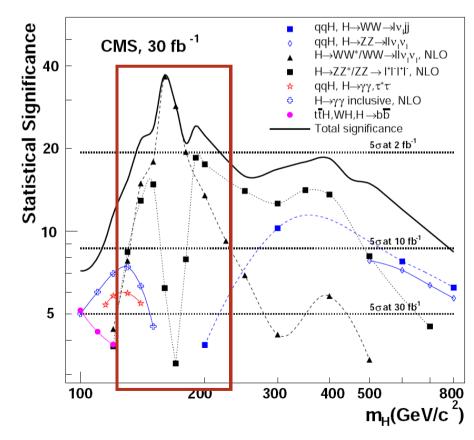
- How to detect Higgs Bosons ?
 - Essential for Higgs discovery is:

[production rate] × [decay probability] × [detection efficiency]

- Higgs events need to be silhouetted against *huge* amount of non-Higgs events \rightarrow e.g. hopeless to see $H \rightarrow b\overline{b}$ via gluon fusion
- * signal significance for Higgs detection @ LHC:

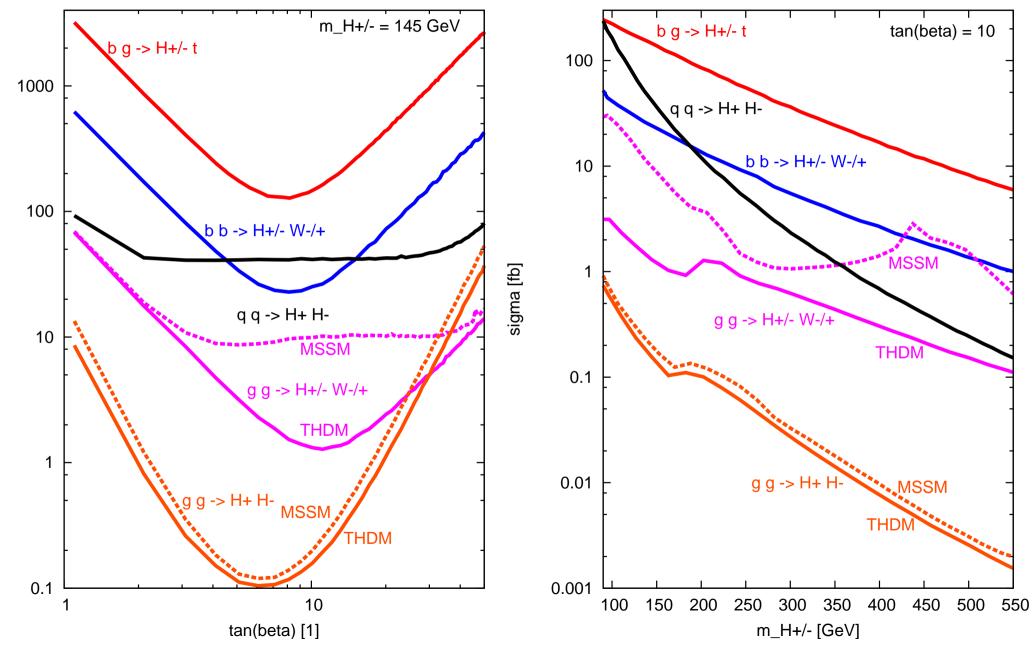
* SM Higgs decay probability (branching ratio):





[Electroweak Interaction, Higgs Bosons & beyond, How to find Higgs Bosons?]

I Predictions: charged Higgs cross sections @ LHC:



[Electroweak Interaction, Higgs Bosons & beyond]

- What else to expect at the LHC?
- Naturalness Problem in the Higgs sector
- Naturalness ['t Hooft 1980]:
- m is a natural small parameter \Longleftrightarrow additional symmetry for $m \rightarrow 0$
- example : electron mass m_e is a natural small parameter:
 - $m_e \rightarrow 0 \implies$ chiral symmetry
 - ullet all quantum corrections to electron self energy $\Sigma_e \propto m_e$
 - partial symmetry protects Σ_e from large quantum corrections
- counter example : SM Higgs mass m_H is not a natural small parameter
 - Higgs Potential: $V_{\text{Higgs}} = -\frac{m_H}{2} \Phi^{\dagger} \Phi + \frac{\lambda}{4} (\Phi^{\dagger} \Phi)^2$
 - $m_H \rightarrow 0 \implies$ no additional symmetry
 - Higgs self energy not protected from large quantum corrections

The Naturalness Problem (also often called "the hierarchy problem"):

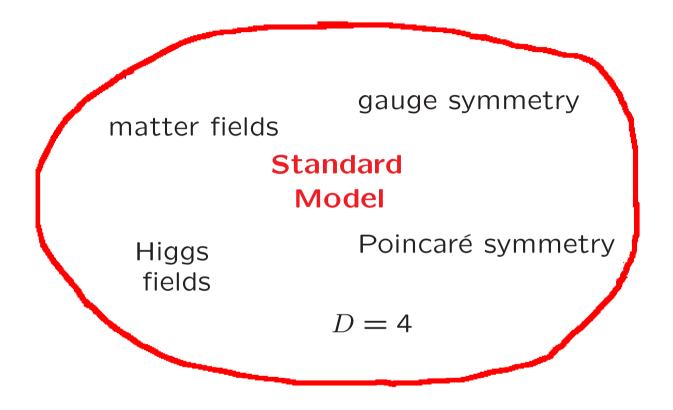
Assuming the SM is only valid up to some scale Λ (say M_{GUT} or M_{Planck}), quantum corrections to the Higgs self energy are of the order of Λ . But present observations indicate a value around the electroweak scale Λ_{FW} .

$$\Lambda_{\rm EW} \propto 100 \, {
m GeV}, \qquad M_{\rm GUT} \propto 10^{15} \, {
m GeV} \qquad M_{\rm Planck} \propto 10^{19} \, {
m GeV}$$

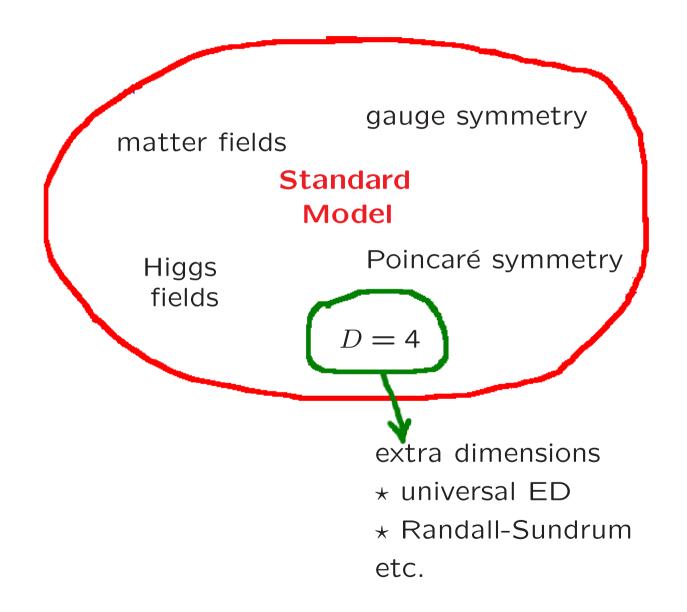
[Electroweak Interaction, Higgs Bosons & beyond, What else to expect at the LHC?]

Taking the Naturalness Problem seriously: What extension of the SM at higher energy scales could avoid the large quantum corrections in the Higgs sector?

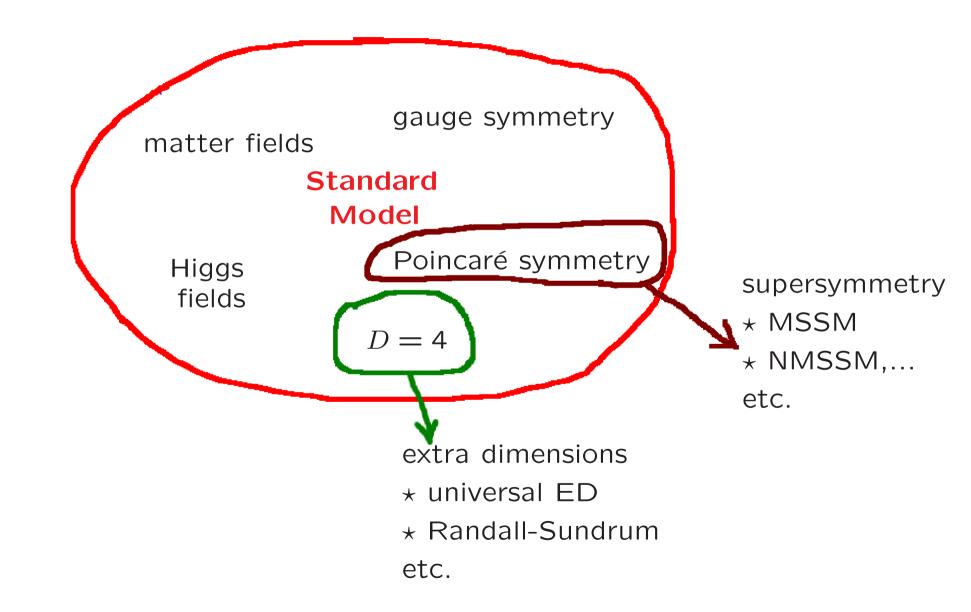
A broad view on SM extensions

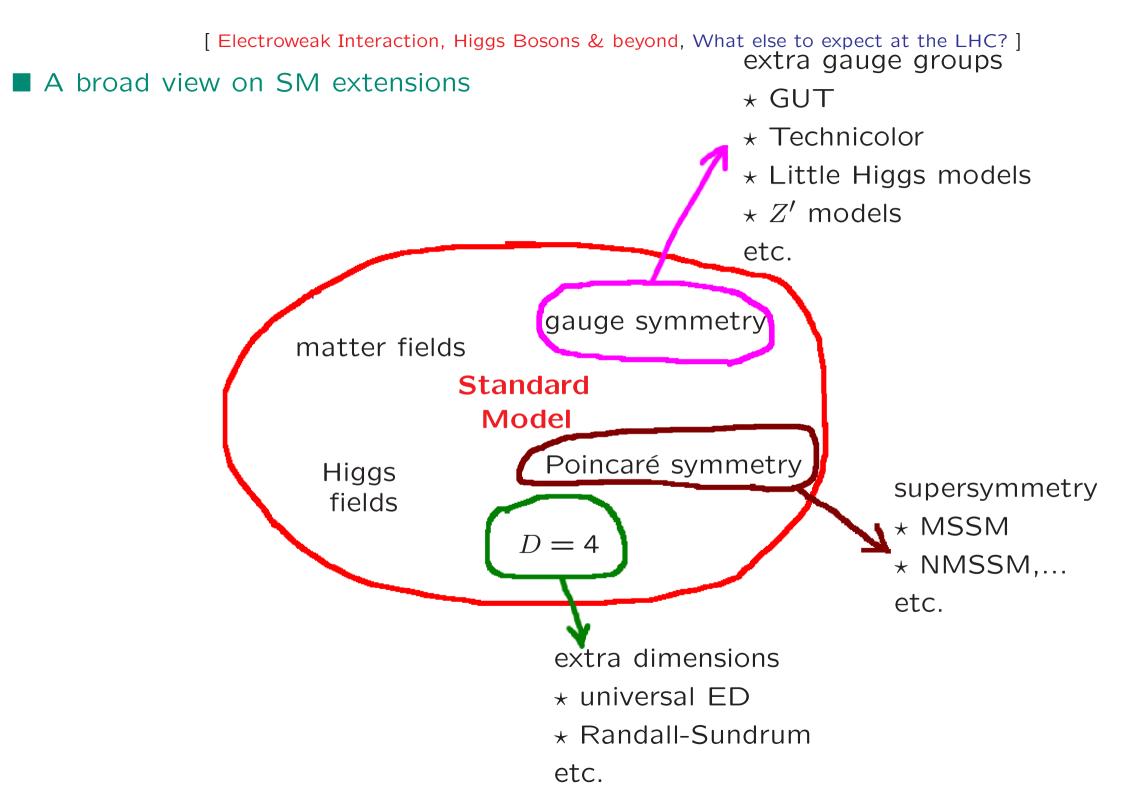


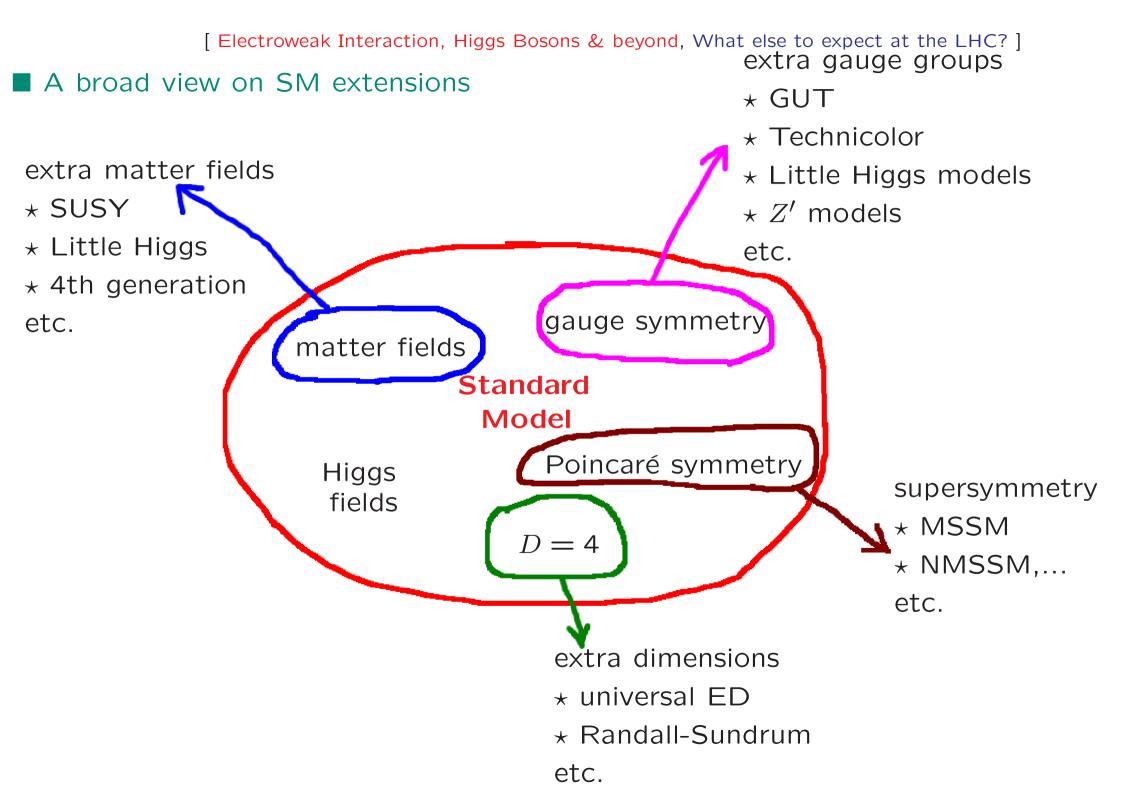
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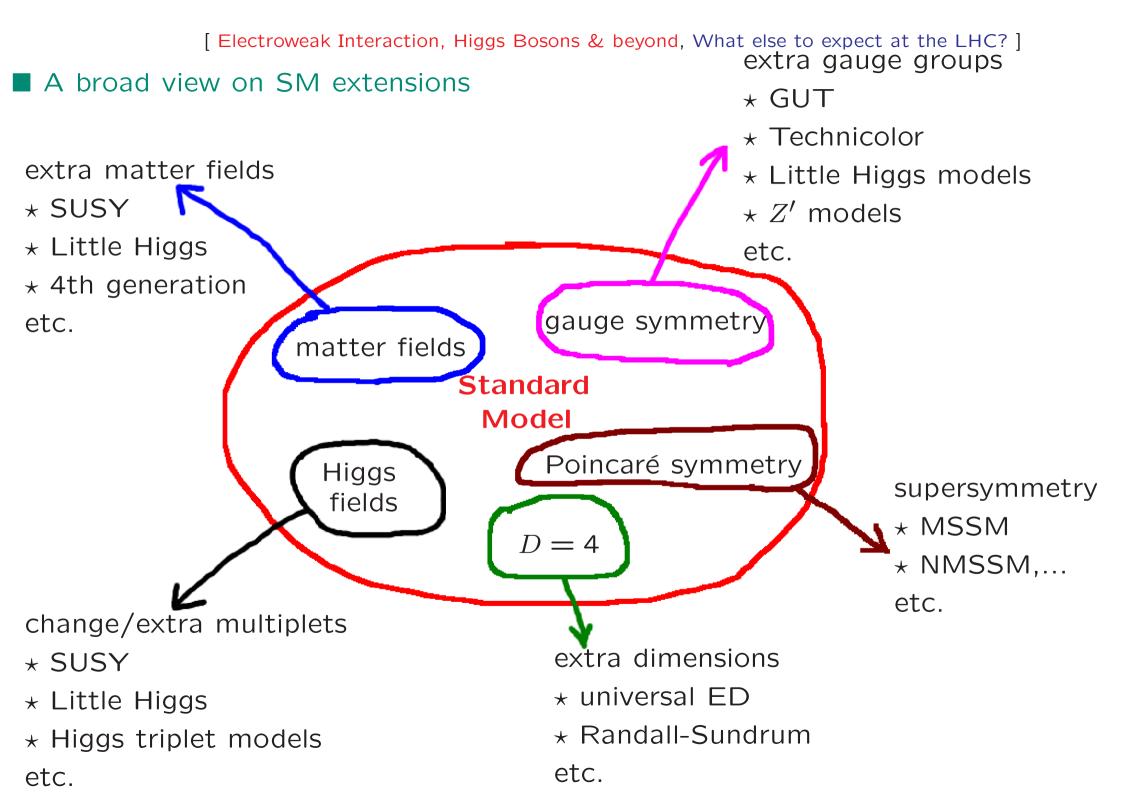


A broad view on SM extensions



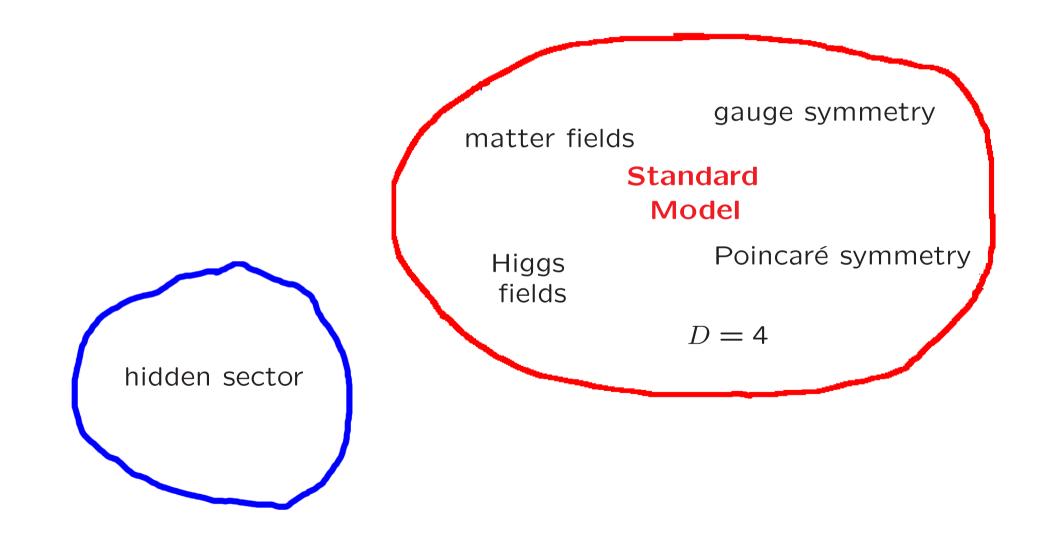






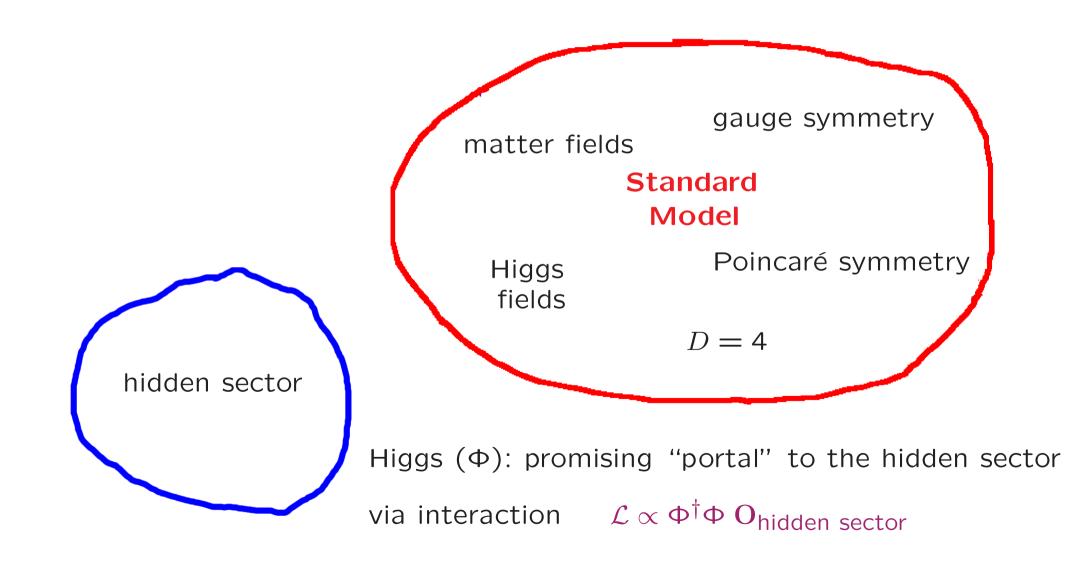
[Electroweak Interaction, Higgs Bosons & beyond, What else to expect at the LHC?]

A broad view on SM extensions



[Electroweak Interaction, Higgs Bosons & beyond, What else to expect at the LHC?]

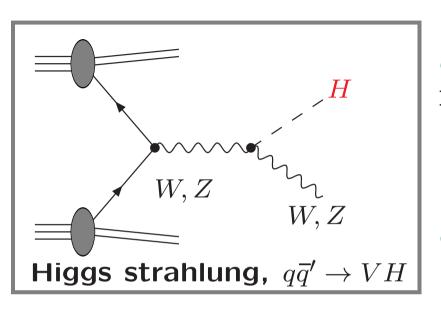
A broad view on SM extensions



• Selected Projects

- SM Higgsstrahlung (NNLO QCD)

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Our calculation: [OBr, Djouadi, Harlander '03] Observation 1: In LO/NLO QCD the cross section factorizes (V = W, Z): $\frac{d\sigma}{dk^2}(q\bar{q} \rightarrow HV) = \sigma(q\bar{q} \rightarrow V^*(k)) \cdot \frac{d\Gamma}{dk^2}(V^*(k) \rightarrow HV)$. Observation 2: Complete NNLO QCD corr. to $\sigma(q\bar{q} \rightarrow V^*)$ are known [Hamberg, van Neerven, Matsuura '91; Harlander, Kilgore '02].

 \rightarrow Idea : Use $\sigma_{NNLO}(q\bar{q} \rightarrow V^{\star})$ to evaluate $\sigma(pp \rightarrow HV)$.

status of theory predictions:

SM, LO [Glashow, Nanopoulos, Yildiz '78]

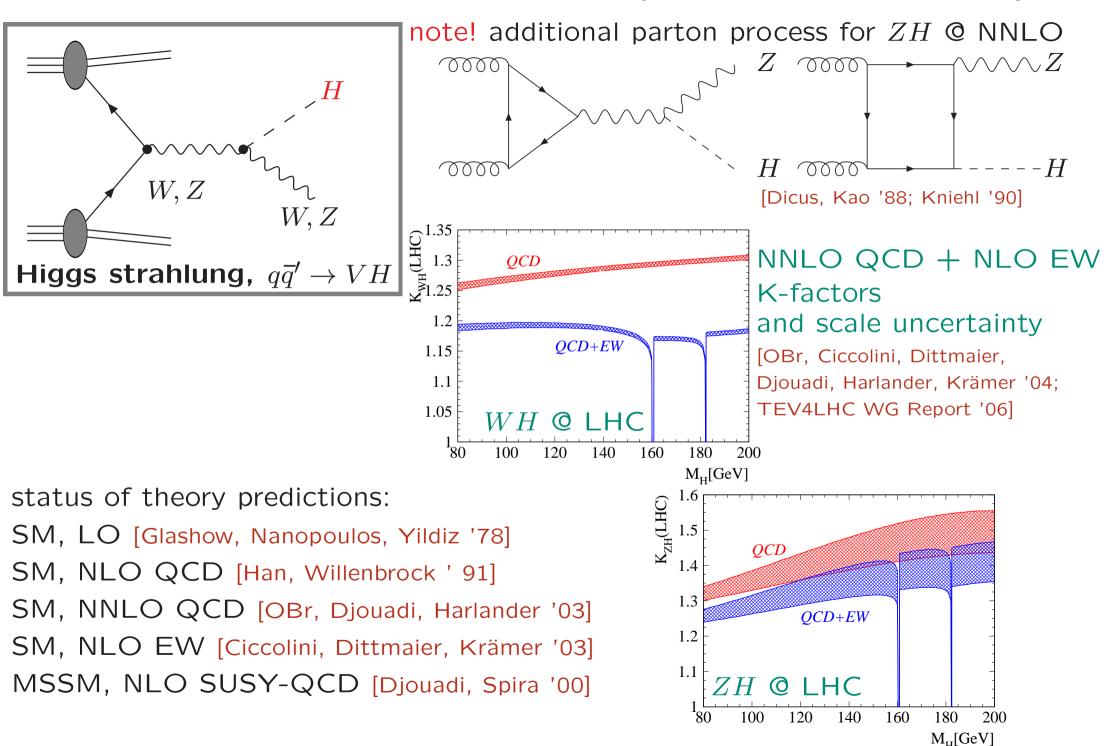
SM, NLO QCD [Han, Willenbrock ' 91]

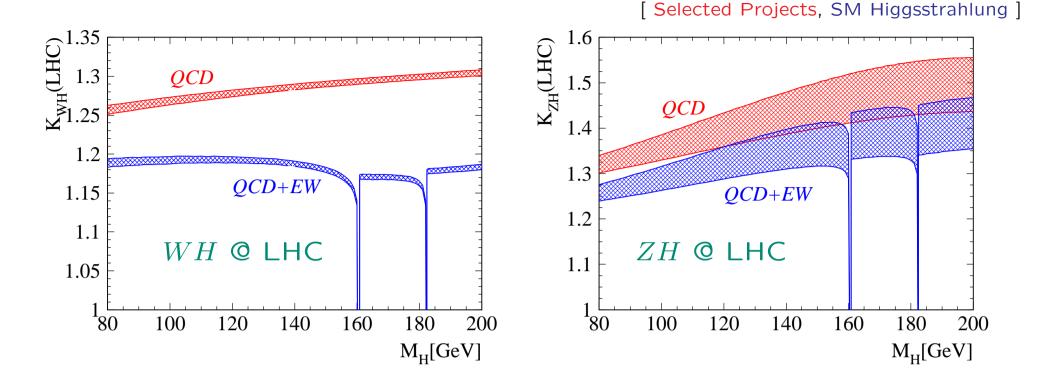
SM, NNLO QCD [OBr, Djouadi, Harlander '03]

SM, NLO EW [Ciccolini, Dittmaier, Krämer '03]

MSSM, NLO SUSY-QCD [Djouadi, Spira '00]

[Selected Projects, SM Higgsstrahlung]





- most precisely known Higgs production process at hadron colliders
- results regularly used by Tevatron collaborations
- recently, we provided updated predictions for total cross sections and uncertainties within the LHC Higgs Cross Section Working Group [CERN Yellow Report 2011]
- \rightarrow ongoing effort, now focusing on differential distributions
- \rightarrow code vh@nnlo to go public soon [OBr, Harlander, Zirke '11]

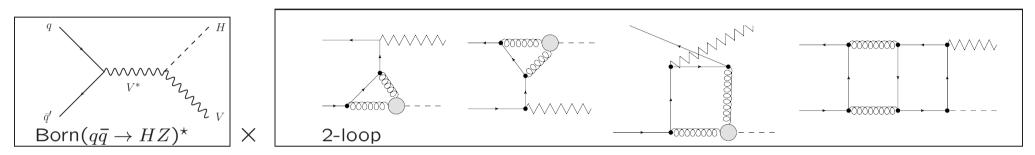
Top quark induced corrections [OBr. Harlander, Wiesemann, Zirke '11] Unfortunatley, this is not the whole story! : Top quark induced corrections appear at NNLO QCD

and are not Drell-Yan like.

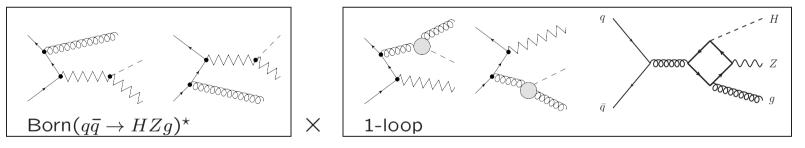
• virtual corrections at NNLO to $q\bar{q} \rightarrow HZ$:

(shaded blob = top quark loop)

 \rightarrow previously overlooked!



• real corrections at NNLO: $(q\bar{q} \rightarrow HZg, qg \rightarrow HZq, \bar{q}g \rightarrow HZ\bar{q})$



• similar corrections for WH production (and in vector boson fusion too!)

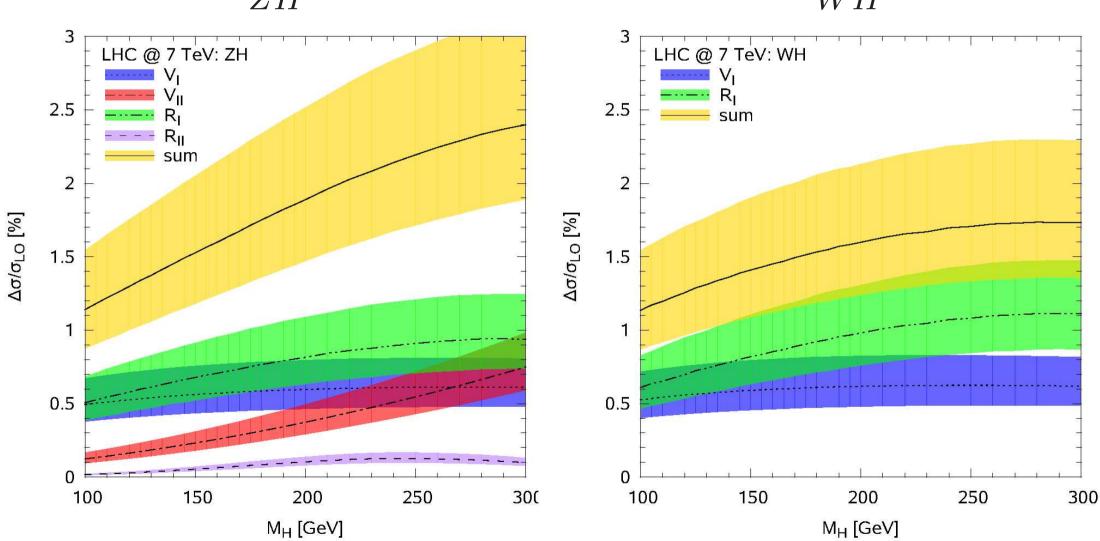
• technical challenges: 1) agreement between two indep. calculations using: a)... asymptotic expansions b)... effective vertices & tensor reduction 2) exact calculation of real corrections

[Selected Projects, SM Higgsstrahlung]

[Selected Projects, SM Higgsstrahlung]

■ LHC(7 TeV) results :

[OBr, Harlander, Wiesemann, Zirke '11]



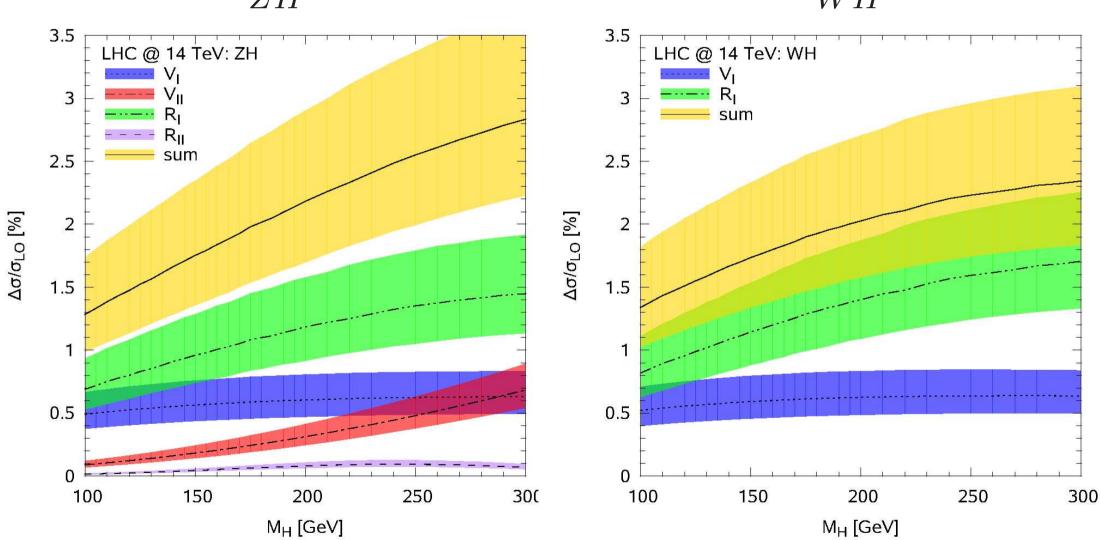
ZH

WH

■ LHC(14 TeV) results :

[Selected Projects, SM Higgsstrahlung]

[OBr, Harlander, Wiesemann, Zirke '11]



ZH

WH

[Selected Projects, SM Higgsstrahlung]

[OBr, Harlander, Wiesemann, Zirke '11]

2.5 2.5 Tevatron: WH Tevatron: ZH V ----- V_I V_{II} R sum R 2 2 R_{II} sum 1.5 1.5 Δσ/σ_{LO} [%] Δσ/σ_{LO} [%] 1 1 0.5 0.5 0 0 100 120 140 160 180 200 100 120 140 160 180 200 M_H [GeV] M_H [GeV]

ZH

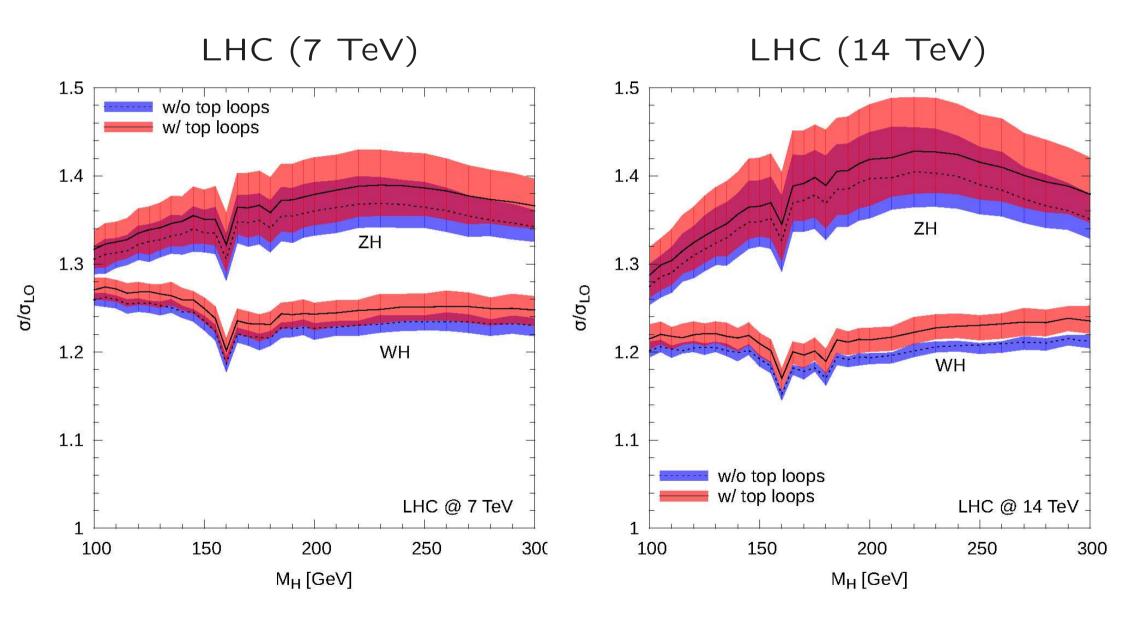
WH

Tevatron results :

[Selected Projects, SM Higgsstrahlung]

LHC results : change in NNLO/LO K-factor

[OBr, Harlander, Wiesemann, Zirke '11]



- Randall-Sundrum scalar sector constrained

- Randall-Sundrum scalar sector constrained
- HiggsBounds : [Bechtle, OBr, Heinemeyer, Stefaniak, Weiglein, Williams '08-'11] tests models with arbitrary Higgs sectors against exclusion bounds from direct searches.
 - easy access to all relevant Higgs exclusion limits including information not available in the publications. (e.g. expected 95% CL cross section limits)
 - applicable to models with arbitrary Higgs sectors (narrow widths assumed) HiggsBounds Input: the predictions of the model for: # of neutral & charged Higgs bosons h_i , m_{h_i} , $\Gamma_{tot}(h_i)$, BR($h_i \rightarrow ...$), production cross section ratios (wrt reference values)
 - combination of results from LEP, Tevatron and LHC possible
 - three ways to use HiggsBounds:
 □ command line, □ subroutines (Fortran [77]/90), □ web interface:
 projects.hepforge.org/higgsbounds

HiggsBounds: status and outlook

- The code is publicly available since Feb. 2009 (current version: 3.5.0 beta) \rightarrow projects.hepforge.org/higgsbounds
 - Tevatron & LHC results up to Lepton-Photon 2011 included
 - extended functionality (H^{\pm} searches, onlyP analyses selection, ...)
 - HiggsBounds 3.5.0 beta available to download
 - \rightarrow includes: SLHA input option, χ^2 fitting for LEP channels, optional addition: SusyBounds (Chargino, Neutralino bounds)
- Reception very good (> 100 users). Code used in/by: FeynHiggs, CPsuperH, Fittino, MasterCode, 2HDMC, DarkSusy, SuperIso, etc.
- Current work/plans:
 - new LHC results after Lepton-Photon 2011
 - searches for fermiophobic models
 - doubly charged Higgs searches, LEP searches for $m_{H} < 10 \ {\rm GeV}$
 - inclusion of width-dependent limits

[Selected Projects, RS scalar sector constrained]

Randall Sundrum model basics:

[Randall, Sundrum '99]

• space has D = 3 + 1 dimensions, metric:

$$ds^{2} = e^{-2kr_{c}\phi}\eta_{\mu\nu}dx^{\mu}dx^{\nu} - r_{c}^{2}d\phi^{2}, \ \phi \in [0,\pi].$$

Spacetime is a slice of 5d anti-de-Sitter space: two boundaries: $\phi = \pi$: IR brane (our 3-space) $\phi = 0$: UV brane

- k, r_c^{-1} are $\mathcal{O}(M_{\mathsf{Pl}})$ with $kr_c \approx 12$. This "little hierarchy" can be generated & stabilized [Goldberger, Wise '00] \Rightarrow fluctuations of r_c : scalar d.o.f φ , gets a VEV Λ_{ϕ}
- resolution of the hierarchy problem: Why is the EW scale $<< M_{Pl}$?: mass parameters in the fundamental 5d model m_0 appear in our visible space as:

$$m = m_0 e^{-kr_c\pi} \approx m_0 10^{-16}$$
.

• propagating in extra dimension:

originally: only gravity,

nowadays: gauge bosons, fermions [EW & flavour observables!] But: Higgs needs to be localized on/near IR brane [hierarchy problem!]

Randall Sundrum scalar sector:

- There is one graviscalar in 5d: the radion φ (typically the lightest new particle to appear)
- Higgs radion mixing via the interaction

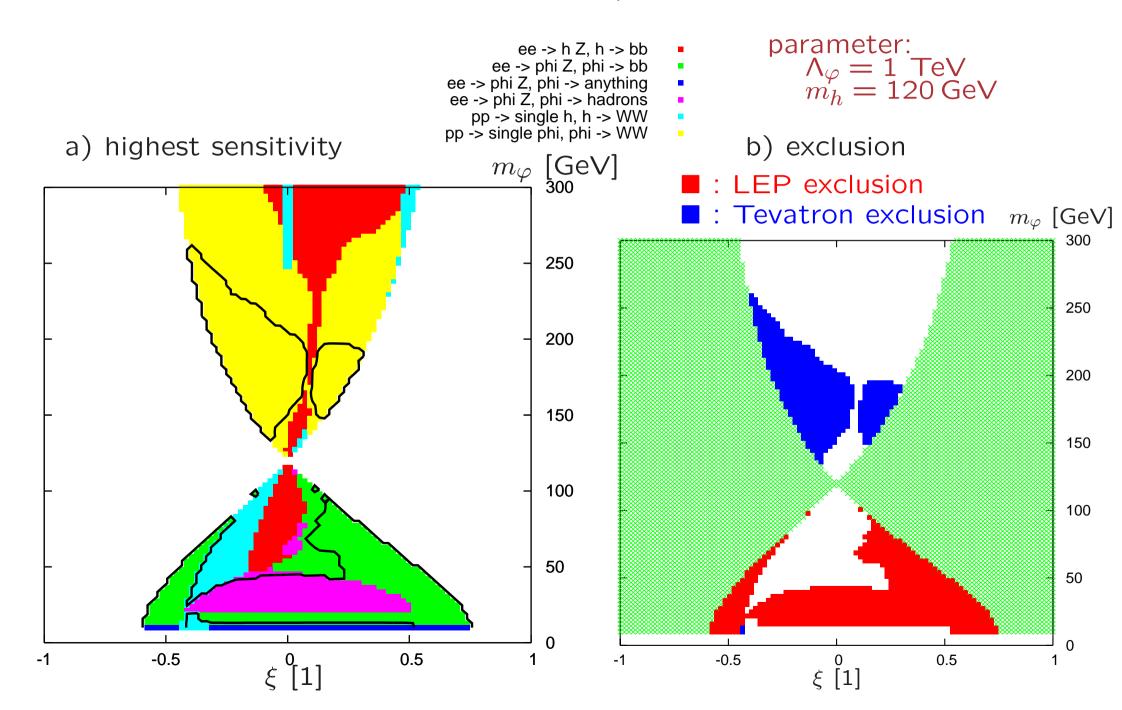
 $\mathcal{L} = -\xi \sqrt{-g_{\text{ind}}} R(g_{\text{ind}}) \Phi^{\dagger} \Phi$

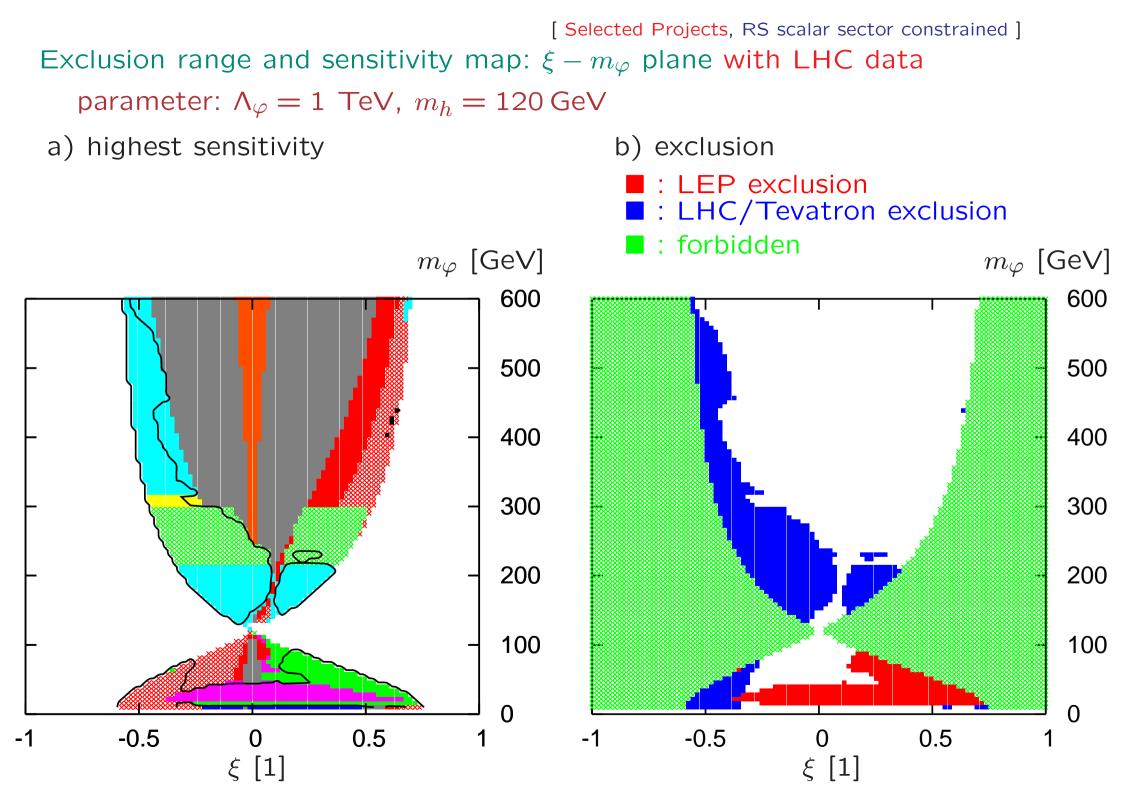
with $g_{ind}(\varphi(x),...)$: induced 4d metric on IR brane, R: Ricci scalar. \rightarrow Radion φ and physical Higgs h mix to form two mass eigenstates

- $-\varphi$ coupling to massive fermions and gauge bosons \propto mass, but
 - $\star \varphi b \overline{b}$ coupling suppressed wrt SM Higgs
 - $\star \varphi gg$ coupling enhanced wrt SM Higgs
 - $\star \varphi \gamma \gamma$ coupling suppressed wrt SM Higgs
- \rightarrow two scalars in the spectrum with modified couplings compared to the SM Higgs boson

[Selected Projects, RS scalar sector constrained]

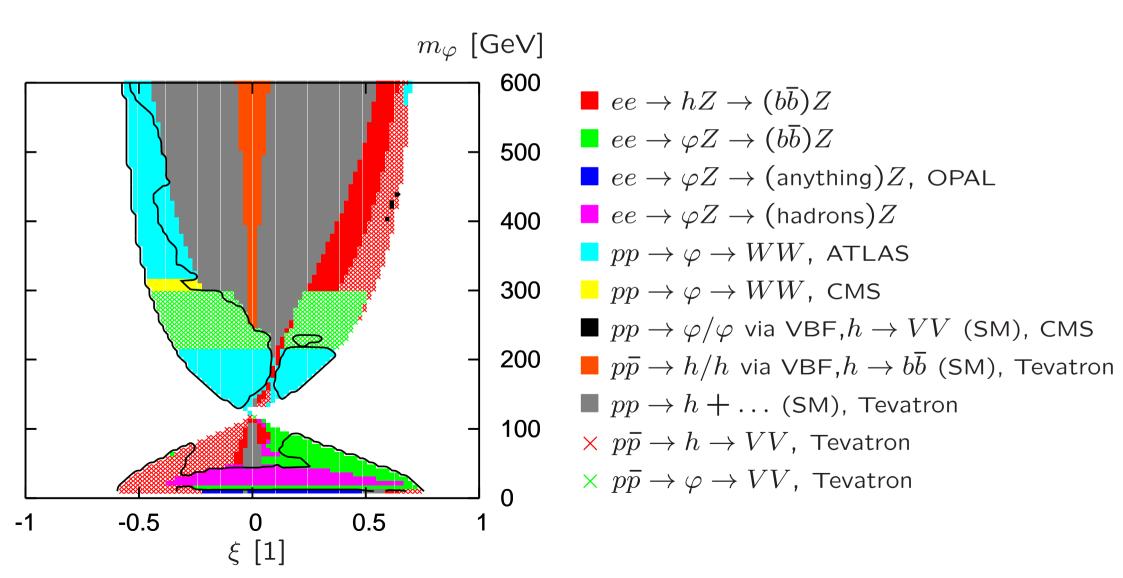
Exclusion range and sensitivity map: $\xi - m_{\varphi}$ plane w/o LHC data (12/2010)





[Selected Projects, RS scalar sector constrained] Exclusion range and sensitivity map: $\xi - m_{\varphi}$ plane with LHC data parameter: $\Lambda_{\varphi} = 1$ TeV, $m_h = 120 \text{ GeV}$

a) highest sensitivity

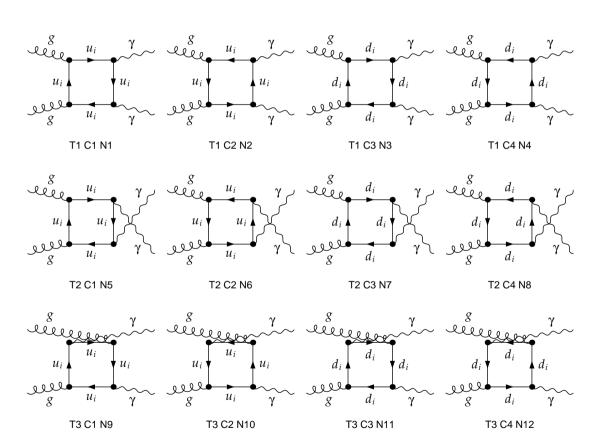


– New physics in $\gamma\gamma/WW/ZZ$ production

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- squark & Kaluza-Klein quark contributions to $gg\to\gamma\gamma$

SM process

 $g \quad g \quad \rightarrow \quad \gamma \quad \gamma$



In general:

all particles carying colour and electrical charge contribute.

Supersymmetry (MSSM):

- additional contributions by quark superpartners (squarks)
- squark masses $\propto M_{\rm SUSY}$

Universal Extra Dimensions (UED):

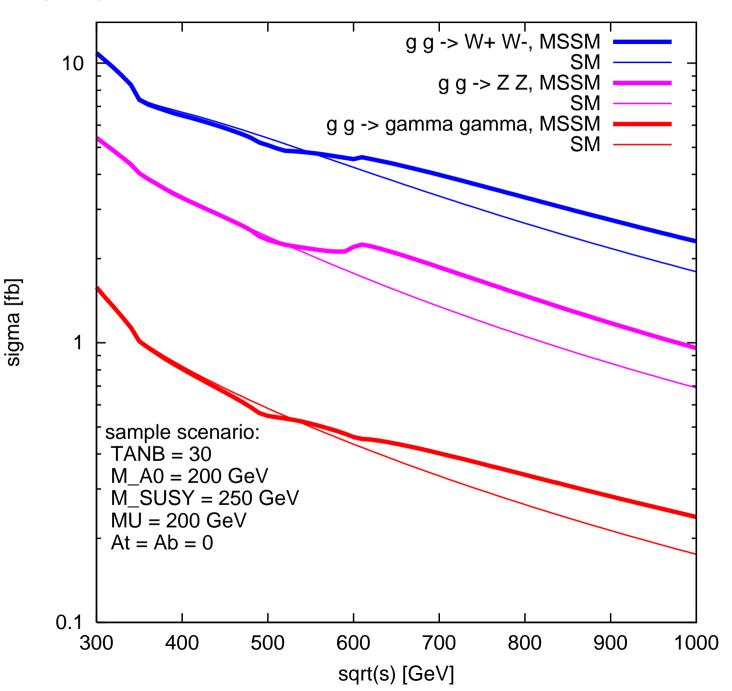
- additional contributions by
 Kaluza-Klein (KK) excitations
 of the quarks
- KK quark masses:

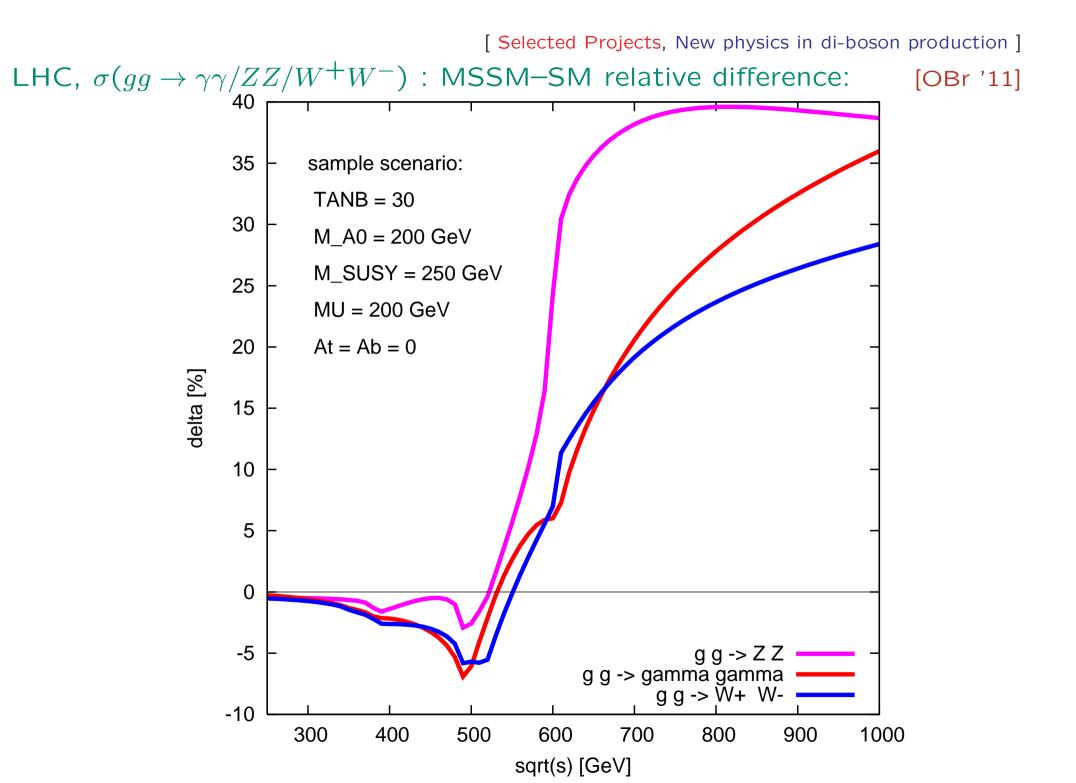
 $m_q^{(n)} = \sqrt{m_q^2 + n^2 m_{\rm KK}^2} \approx n \, m_{\rm KK}$

[Selected Projects, New physics in di-boson production]

[OBr '11]

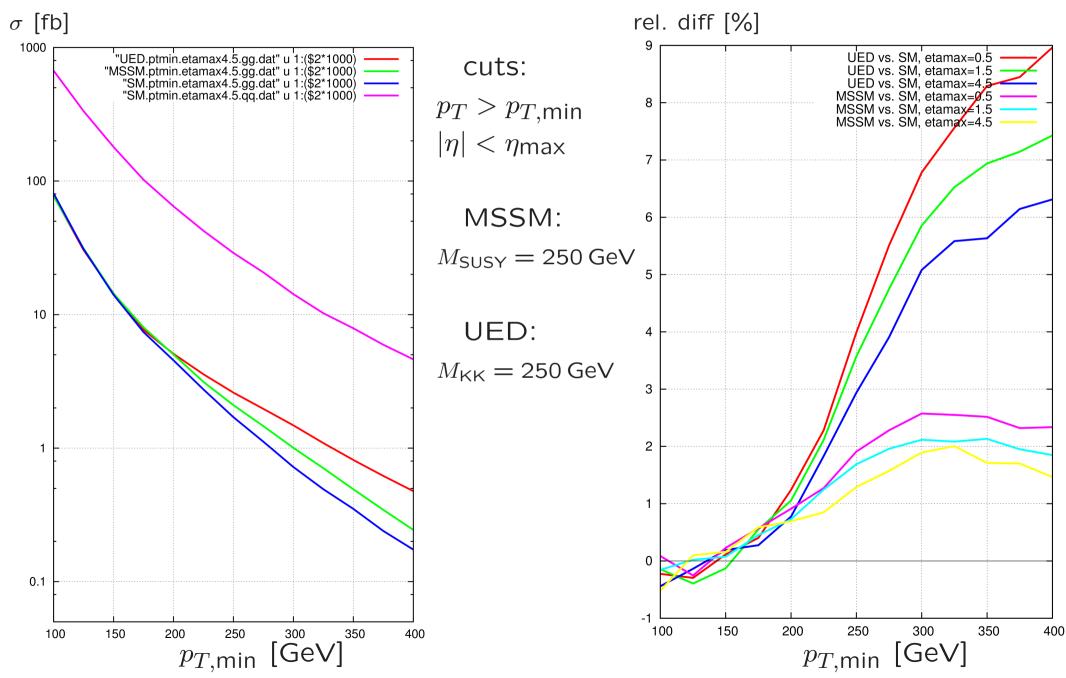
LHC, $\sigma(gg \rightarrow \gamma \gamma/ZZ/W^+W^-)$ in the MSSM:





[Selected Projects, New physics in di-boson production]

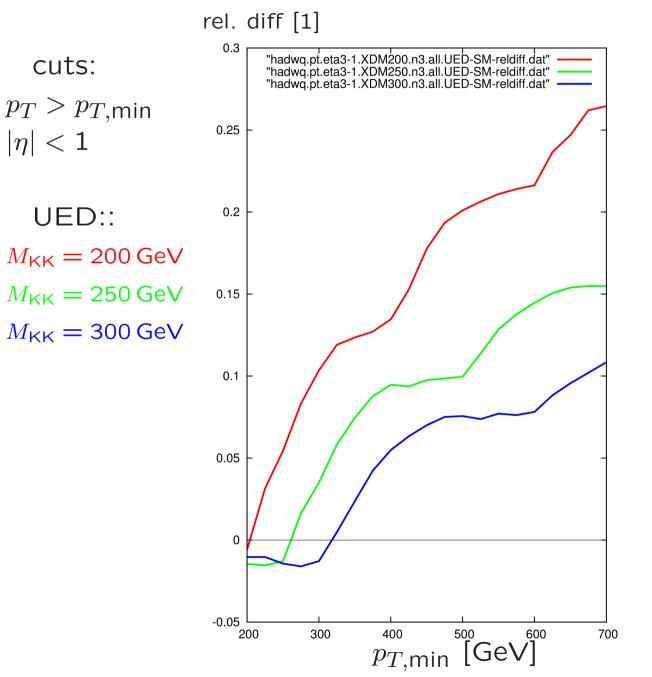
LHC(14 TeV), $\sigma(pp \rightarrow \gamma\gamma)$: UED/MSSM–SM relative difference: [OBr '11]



[Selected Projects, New physics in di-boson production]

[OBr '11]

LHC(14 TeV), $\sigma(pp \rightarrow \gamma\gamma)$: UED–SM relative difference:



summary

- We are sure to observe electroweak symmetry breaking in nature. However, up to now, we have no clue how it is realised. The Higgs mechanism allows to describe EWSB consistently up to very high energy. → Maybe we know more tomorrow!
- SM Higgsstrahlung is now really known at NNLO QCD accuracy. The impact of top-induced corrections is small (adding at most 14 % to the previously known NNLO corrections). Still, SM predictions for vector boson fusion should also be revisited.
- LHC/Tevatron Higgs searches rule out new parts of the Randall-Sundrum model's parameter space (compared to LEP results).
- The di-photon and WW/ZZ production show potential for the discrimination between models. Further investigations are needed.

• Backup

- MSSM

Supersymmetry ...

- ... is *the* extension of the Poincaré-symmetry of space-time
- ... leads to a symmetry between Fermions & Bosons
- gauge theory with minimal SUSY :
 - same # of fermionic & bosonic d. o. f.
 - \rightarrow a superpartner of different spin exists for each particle
 - couplings are correlated
 - \rightarrow e.g. scalar 4-point int. \leftrightarrow gauge couplings
 - superpartners have the same mass
 - \rightarrow SUSY must be broken at the electroweak scale
- gauge theory with broken SUSY :
 - superpartner masses enter as additional free parameters (essentially)

Minimal supersymmetric Standard Model (MSSM):

gauge group : $SU(3)_{colour} \times SU(2)_{isospin} \times U(1)_{hypercharge}$

particle content :

regular particles		spin	superpartners		spin
fermions <	$\left\{egin{array}{l} { ext{quarks}}\ u,d,s,c,b,t\ { ext{leptons}}\ e, u_e,\mu, u_\mu, au, u_ au, u_ au ight.$	<u>1</u> 2	sfermions <	$\left\{egin{array}{l} { m squarks}\ { ilde u}, { ilde d}, { ilde s}, { ilde c}, { ilde b}, { ilde t}\ { m sleptons}\ { ilde e}, { ilde u}_e, { ilde \mu}, { ilde u}_\mu, { ilde au}, { ilde u}_ au \end{array} ight.$	0
gauge bosons G, W^{\pm}, Z, γ		1	gauginos	$ ilde{G}, ilde{W}^{\pm}, ilde{Z}, ilde{\gamma}$	$\frac{1}{2}$
Higgs bosons H_1, H_2		0	Higgsinos	$ ilde{H}_1, ilde{H}_2$	$\frac{1}{2}$

 $\tilde{W}^{\pm}, \tilde{Z}, \tilde{\gamma}$ and \tilde{H}_1, \tilde{H}_2 mix to charginos $\chi_1^{\pm}, \chi_2^{\pm}$ and neutralinos $\chi_1^0, \ldots, \chi_4^0$

R-parity : discrete, multiplicative quantum number

R(regular particles) = +1R(superpartners) = -1

 \rightarrow designed to avoid large Flavour Canging Neutral Currents (FCNC)

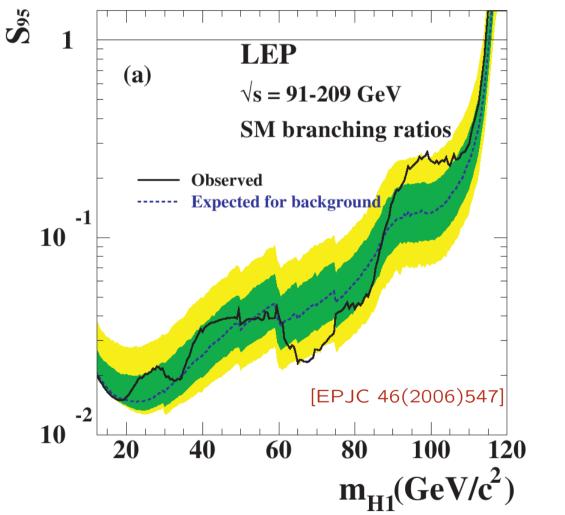
consequences of *R*-parity conservation:

all interactions involve an *even* number of superpartners
 → superpartners can only be pair-produced

• the lightest superpartner (LSP) is stable \rightarrow the LSP is a candidate for dark matter

– HiggsBounds implementation

Higgs search results: example 1: LEP SM combined limit

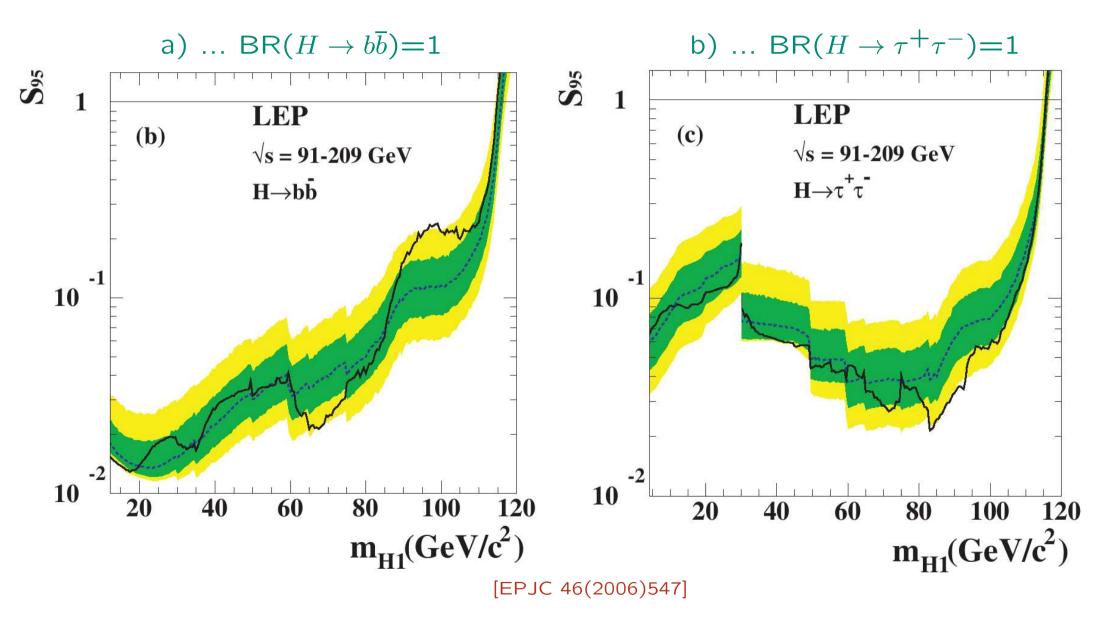


exclusion = rejection of the Higgs hypothesis

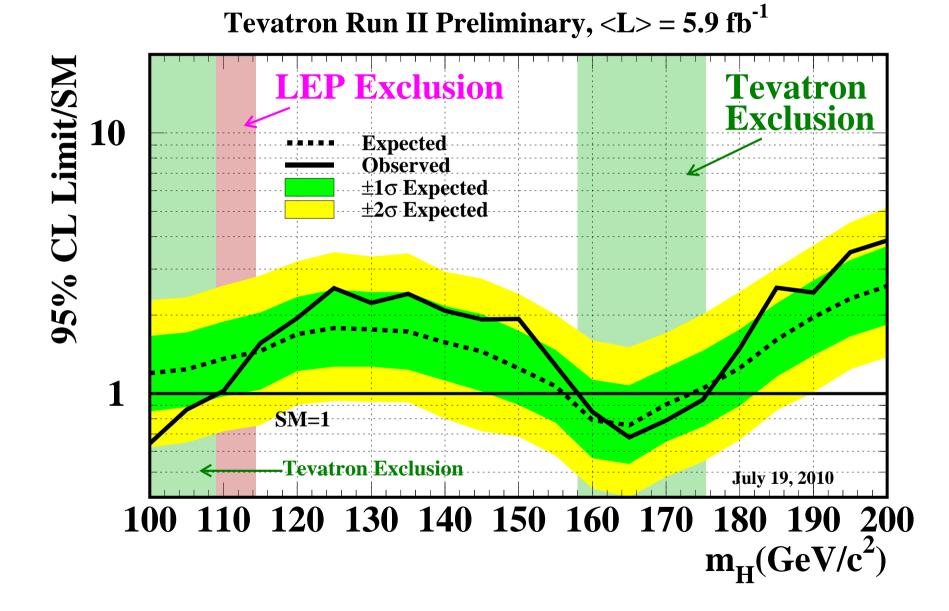
 $S_{95}(m_{H1}) := \frac{\sigma_{\min}}{\sigma_{SM}}(m_{H1})$

where $\sigma_{\min}(m_{H1})$ is the Higgs signal cross section where data and Higgs hypothesis are compatible with only 5% probability.

A SM-like model with $\sigma_{model}(m_{H1}) > \sigma_{min}(m_{H1})$ or $\frac{\sigma_{model}(m_{H1})}{\sigma_{min}(m_{H1})} > 1$ is said to be excluded at the 95% C.L. example 2: LEP single topology limits, assuming HZ production and ...



example 3: Tevatron SM combined limit [CDF & DØ '10]



Considering many analyses for many Higgs bosons:

first a definition : **analysis application** *X*:

application of a certain analysis A_i to a certain Higgs boson h_k (or a set)

that means: X corresponds to:

- * a signal topology (or a set),
- * the corresponding cross section prediction $\bar{\sigma}_{model}(X)$,
- * observed cross section limit $\overline{\sigma}_{\text{observed}}(X)$ of analysis A,
- * expected cross section limit $\bar{\sigma}_{\text{expected}}(X)$ of analysis A.

Basic idea:

for an analysis application X:

• evaluate model prediction

$$\bar{\sigma}_{\text{model}}(X) = \frac{[\sigma \times BR]_{\text{model}}}{[\sigma \times BR]_{\text{ref}}}$$
 (reference: usually SM

of the correponding search topology for given Higgs masses + deviations from the reference.

- read off the corresponding observed 95% C.L. limit: $\overline{\sigma}_{\text{observed}}(X)$.
- If $\overline{\sigma}_{model}(X) > \overline{\sigma}_{observed}(X)$ the model is excluded by this analysis application at 95% C.L.
- \rightarrow Problem : how to combine analysis applications without losing the 95% C.L. ?

Answer: We can't do that.

Only a dedicated experimental analysis can do that.

However: we can always use the analysis application of highest statistical sensitivity.

How to preserve the 95% C.L. limit:

- Obtain for each X the experimental expected limit $\bar{\sigma}_{\text{expected}}(X)$.
- Determine the analysis application X_0 with the highest sensitivity for the signal, i.e. of all X, find X_0 where $\frac{\overline{\sigma}_{model}(X)}{\overline{\sigma}_{expected}(X)}$ is maximal.
- If for this analysis application $\overline{\sigma}_{model}(X_0) > \overline{\sigma}_{observed}(X_0)$, the model is excluded at 95% C.L. by X_0 .